



Feasibility study of open charm elliptic flow in CBM



Sélim SEDDIKI ^{1,2,3} and Fouad RAMI ²



Physics motivation

Goal and strategy of the simulations

Reconstruction of the reaction plane

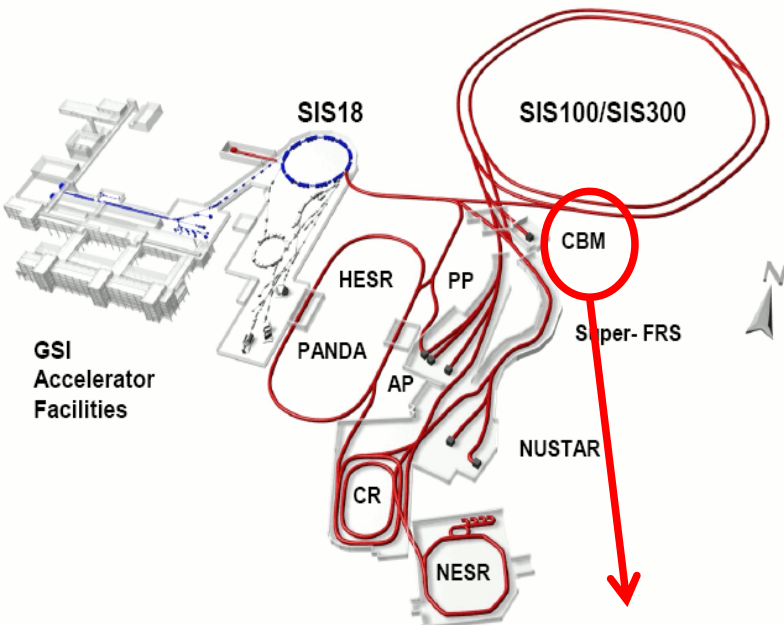
**Feasibility of v_2 measurements for D^+ mesons
(only statistical errors)**

Summary and conclusion



- (1) Institut für Kernphysik and Goethe University, Frankfurt am Main
- (2) Institut Pluridisciplinaire Hubert Curien, Strasbourg
- (3) Helmholtz Research School, Frankfurt am Main

The CBM experiment at FAIR

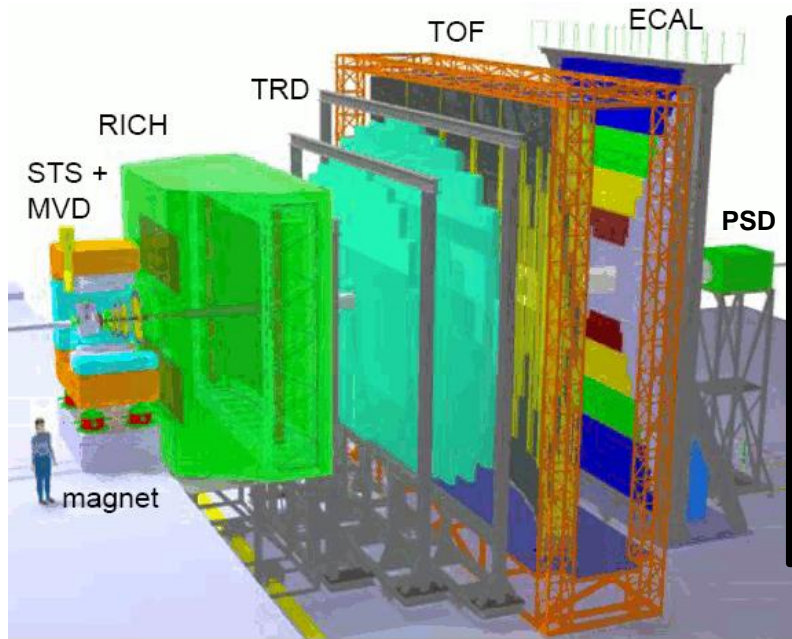


Facility for **A**nti-proton and **I**on **R**esearch
(**FAIR**) at GSI-Darmstadt:

→ **extremely high intensity** HI beam
synchrotron **SIS100/300**

→ **ex: up to 10^{8-9} Au / s @ 45 A.GeV**

CBM: one of the major experiments



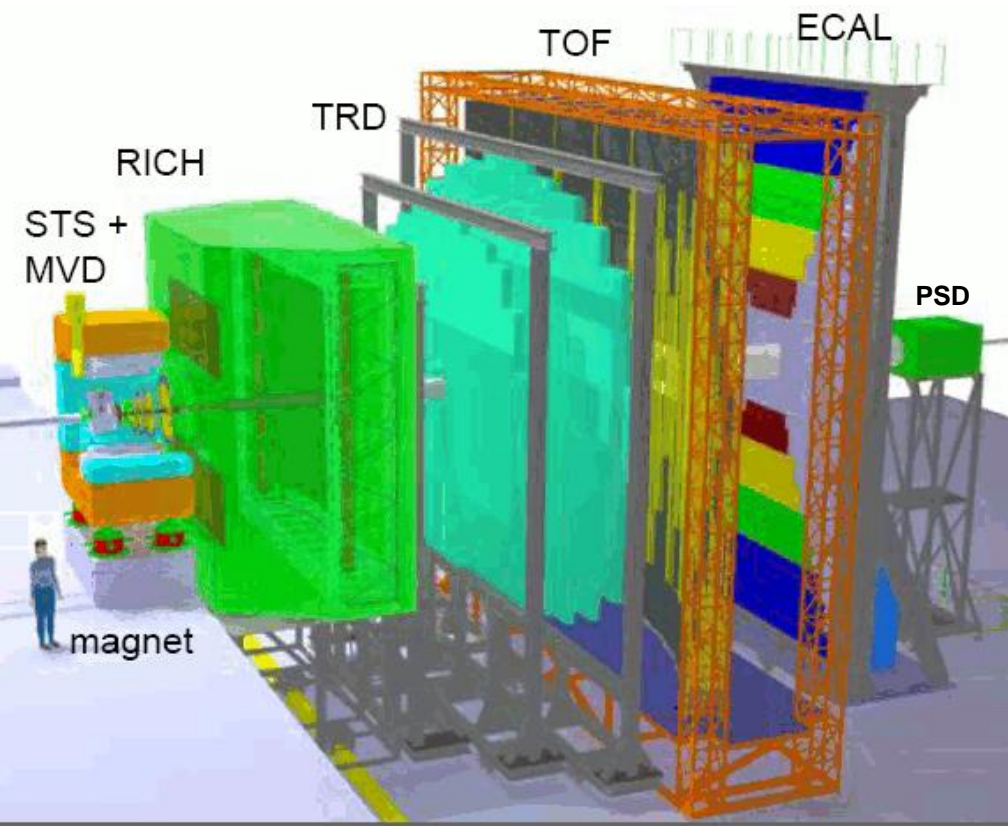
Compressed **B**aryonic **M**atter experiment (CBM):

→ Explore the **nuclear matter phase diagram**
at high baryonic density

→ **Experimental challenge** : measure **rare probes**
- Ξ , Ω , low-mass ρ , ω , ϕ -mesons,
- J/Ψ , Ψ' , open charmed particles
near the threshold energy ($E_c = 15 \text{ AGeV}$)

→ **Extremely fast and radiation-hard** experiment

The CBM experiment at FAIR (2)



Track reconstruction and p-determination:
Silicon Tracking System (STS)

Hadron ID:
STS + Time Of Flight (TOF)

Vertex reconstruction and open charm ID:
Micro-Vertex Detector (MVD)
→ close to interaction point
→ operate in vacuum

Event characterisation :
Projectile Spectator Detector (PSD)

Electron ID / pion suppression:
RICH, TRD, ECAL

Photon detection:
ECAL

Muon detection:
RICH → muon set-up with absorbers

A rich panel of observables foreseen for CBM

↔ Cross-check with independent observables

Deconfinement phase

transition at high ρ_B

- × excitation function and flow of strangeness ($K, \Lambda, \Sigma, \Xi, \Omega$)
- × excitation function and flow of charm ($J/\psi, \psi', D_0, D^\pm, \Lambda_c$)
- × melting of J/ψ and ψ'

QCD critical endpoint

- × excitation function of event-by-event fluctuations ($K/\pi, \dots$)

The equation-of-state at high ρ_B

- × collective flow of hadrons
- × particle production at threshold energies (open charm?)

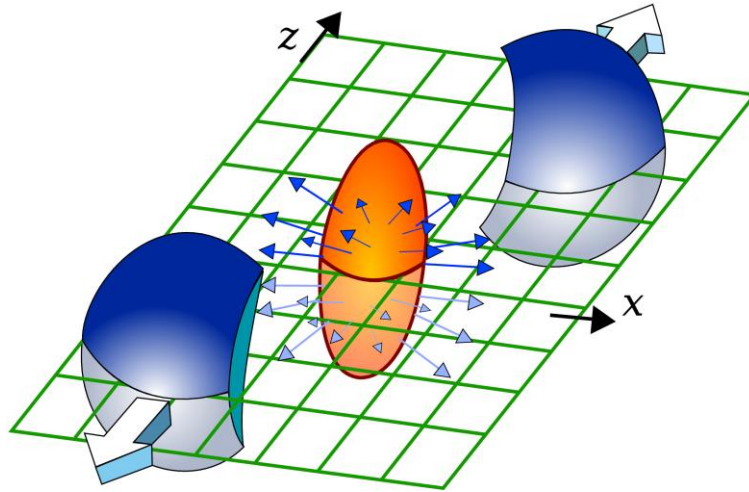
Onset of chiral symmetry restoration at high ρ_B

- × in-medium modifications of hadrons ($\rho, \omega, \phi \rightarrow e^+e^-(\mu^+\mu^-), D$)

- Excitation functions of bulk and rare observables!
- Bulk observables with “unlimited” statistics
- Systematic studies of rare observables (charm, dileptons) with excellent statistics

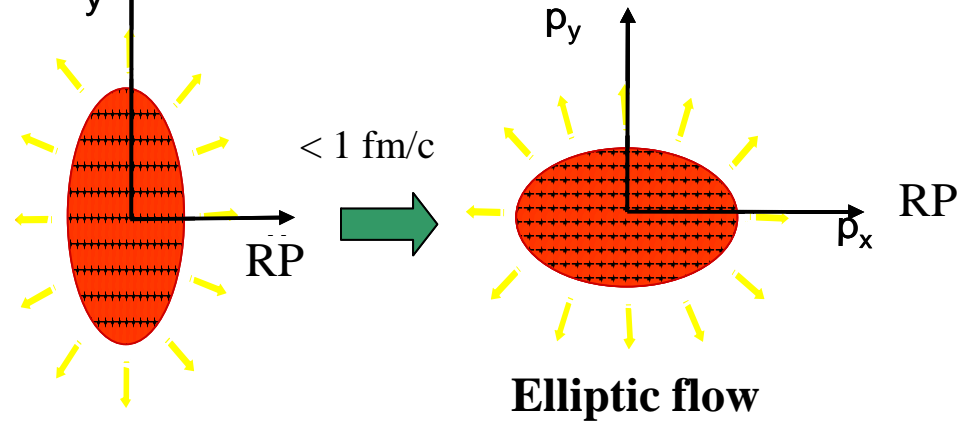
Just a reminder about elliptic flow ...

Reaction plane (RP)

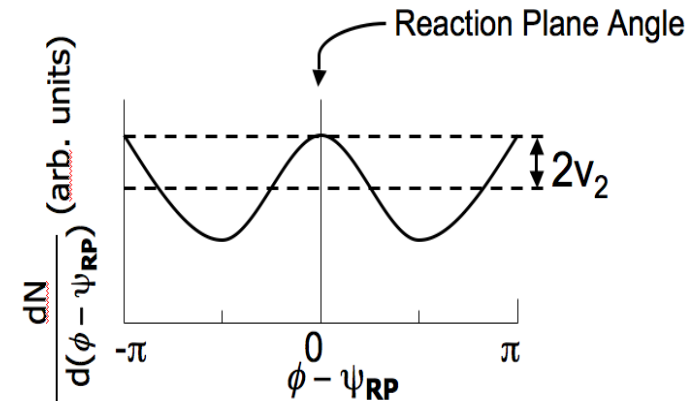


$$e = \frac{\langle y^2 \rangle - \langle x^2 \rangle}{\langle y^2 \rangle + \langle x^2 \rangle}$$

$$v_2 = \frac{\langle p_x^2 \rangle - \langle p_y^2 \rangle}{\langle p_x^2 \rangle + \langle p_y^2 \rangle} = \langle \cos(2 \times (\phi - \Psi_{RP})) \rangle$$



- **I**nitial spatial anisotropy of participants in transverse plane for non-central collisions
 - **V**ia the re-scattering processes in the participant region, pressure gradients appear, higher along the reaction plane, which leads to ...
 - **A** collective flow, preferably « in-plane » (no “squeeze-out” for $E_{\text{beam}} > 5 \text{ A.GeV}$)
- **S**ensitive to the nuclear matter **compressibility** (soft/hard EOS?)
 its **transport properties** (speed of sound, viscosity, etc)
 its **effective dof** (QGP and HRG have \neq transport properties)



$$\frac{dN}{d\Delta\Theta} = \frac{1}{2\pi} \sum_n (1 + 2v_n \cos n\Delta\Theta)$$

Elliptic flow as a probe of the deconfinement phase transition at FAIR energies

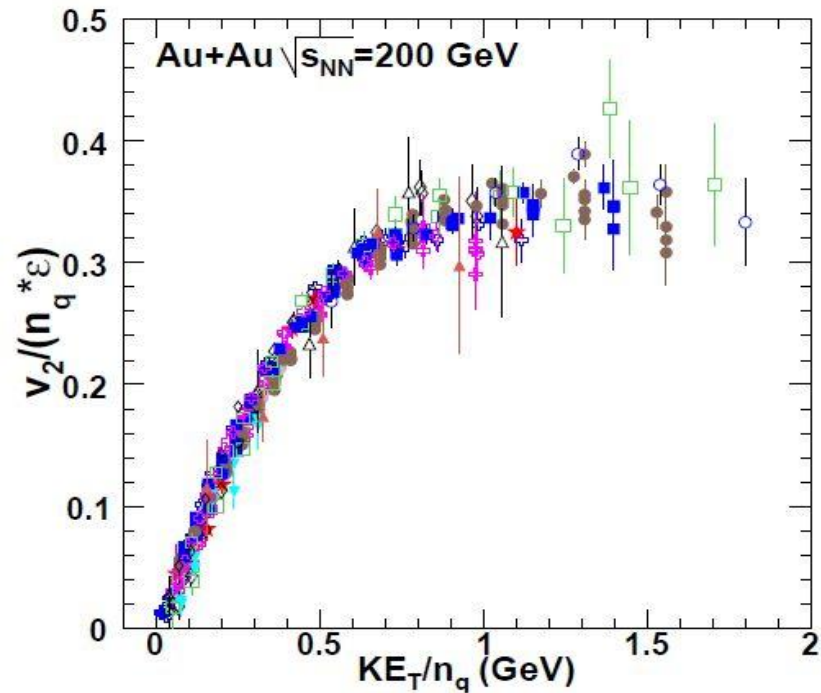
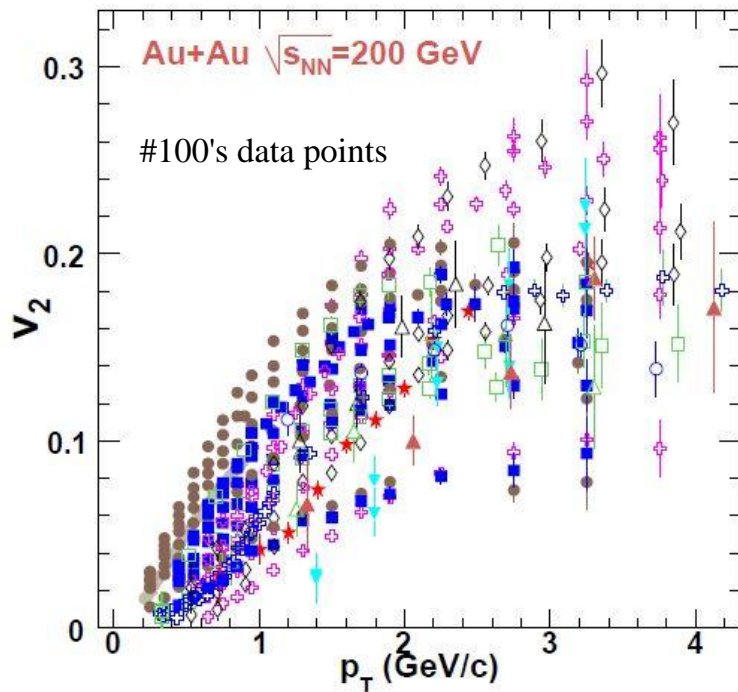
At FAIR energy regime, a **1st order phase transition** is predicted by LQCD to occur in HI collision (for $20 < E_{\text{beam}} < 45$ A.GeV ?):

- predicted collapse of proton v_2 around mid-rapidity at beam energy $E_{\text{lab}} \sim 40$ A.GeV (Stoecker et al., CPOD07_025)
 - indication of this collapse from NA49 data at 40 A.GeV
 - extrapolation from AGS data leads to a collapse of proton v_1 around same energy
 - predicted collapse of proton v_1 at 10 A.GeV by hydro + 1st phase transition (only!)
 - disappearance of the number of valence quark scaling of v_2 (and localisation of the QCD critical point via an energy scan)
 - “large/small” open-charm v_2 ?
- **all these signals require a high precision of the v_2 measurement!**

Disappearance of v_2 CQN scaling as a probe of the phase transition at FAIR energies

- The **Constituent Quark Number (CQN) scaling** of the elliptic flow is one of the strong indications for the **formation of a QGP (RHIC)**

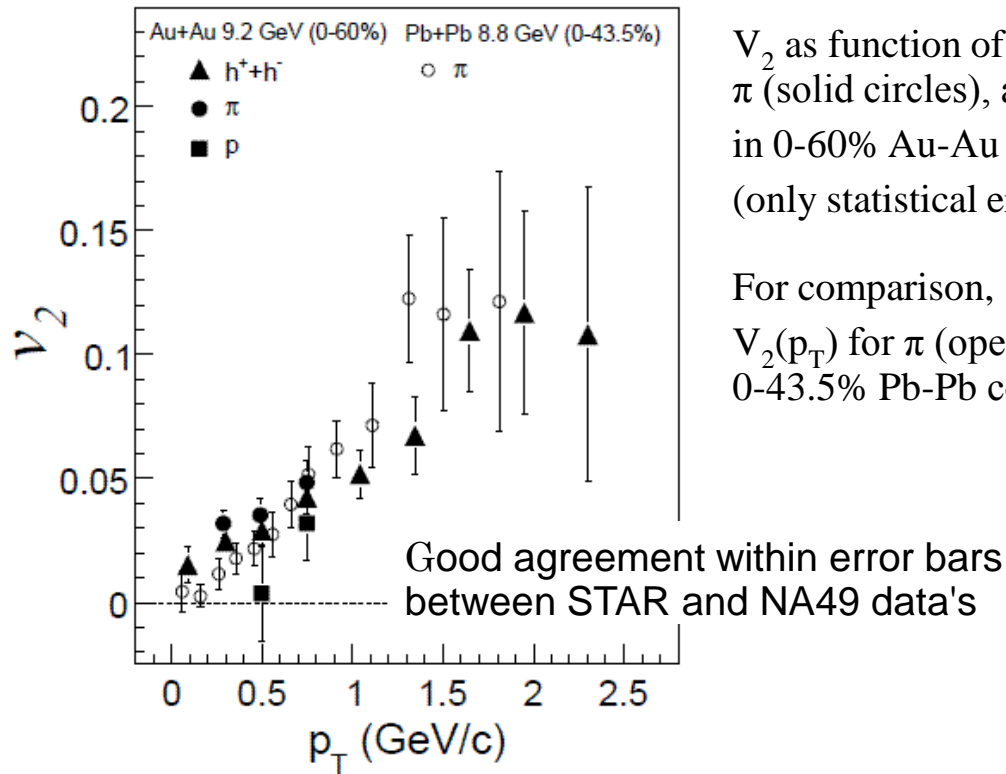
RHIC data, R.Lacey and A.Taranenko, arXiv:nucl-ex/0610029v3



- Search for the **disappearance of this scaling at FAIR energies**
 - 1st order phase transition \leftrightarrow sharp signal
 - energy scan \rightarrow constrain on the localisation of the QCD critical point

Bulk v_2 existing measurements at FAIR energies

ArXiv: 0909.4131v1



V_2 as function of p_T for charged hadrons (solid triangles), π (solid circles), and p (solid squares) in 0-60% Au-Au collisions at $\sqrt{s} = 9.2$ GeV (STAR). (only statistical errors are shown)

For comparison,

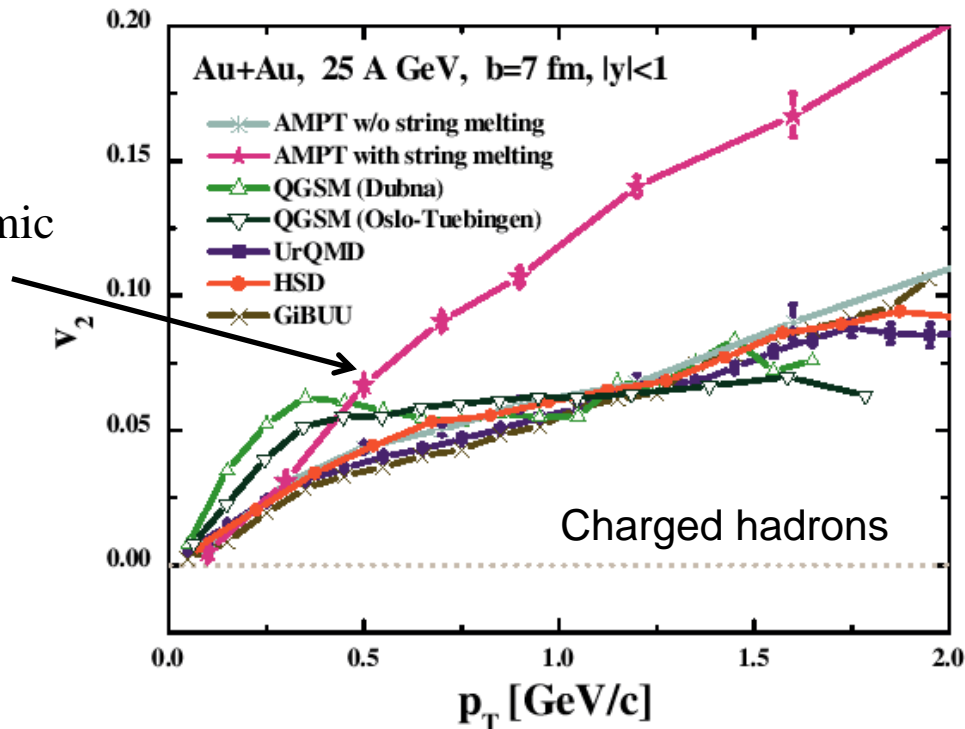
$V_2(p_T)$ for π (open circles) from NA49 in 0-43.5% Pb-Pb collisions at $\sqrt{s} = 8.8$ GeV.

- Existing measurements at RHIC (low energy scan) and SPS, close to FAIR energy regime $\sqrt{s} \sim 7$ GeV
- But insufficient statistics to study the scaling behaviour of the collective flow of particles at $\sqrt{s} \sim 9.2$ GeV (recent data taking – 2010 with higher statistics!)
- The FAIR intense beam intensity will allow for very high statistics studies
- It will give access to the collective flow of rare probes at these energies, namely open charmed particles

Open charm v_2 as a probe of the phase transition at FAIR energies

- The deconfinement phase transition should be visible in the **magnitude of the elliptic flow**

The string melting mimics the partonic dof.



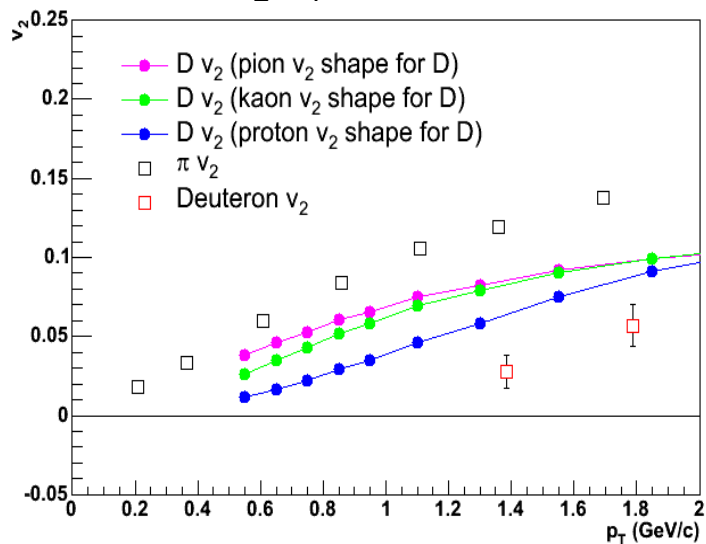
CBM physics book
Submitted to Lecture
Notes in Physic

- The observed v_2 may be bigger in case an early partonic medium contributes to the rescattering process
- **Same argument for open charm v_2** (no prediction so far at FAIR energies):
 - its magnitude is particularly sensitive to the formation of a QGP
 - one of the most challenging differential analysis at FAIR energy regime!

Existing open charm v_2 measurements

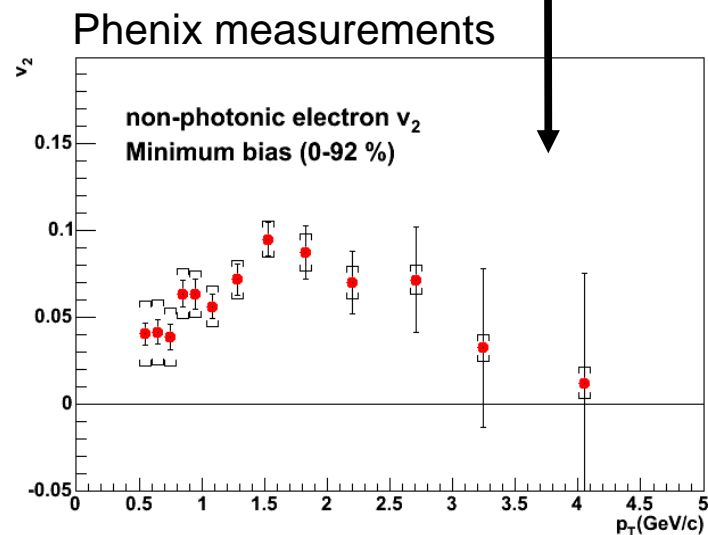
- The open charm v_2 has only been measured via the indirect method using non-photonic electrons:
 - the estimation may suffer from large systematic errors and is model dependant
- A direct measurement of open charm v_2 is one of the motivations for the future Heavy Flavor Tracker (STAR) and Micro-Vertex Detector (CBM)

D $V_2(p_T)$ simulation



Shingo Sakai (PHENIX) (J. Phys G 32, S 551)

- Different assumptions for the shape of D meson $V_2(p_T)$: π, K and p $v_2(p_T)$
- All non-photonic electrons ($p_T < 2.0$ GeV/c) were assumed to come from the D decay
- D \rightarrow e, p_T spectrum constrained by the data

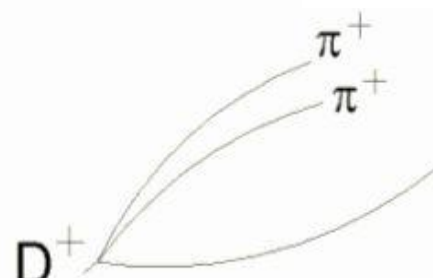
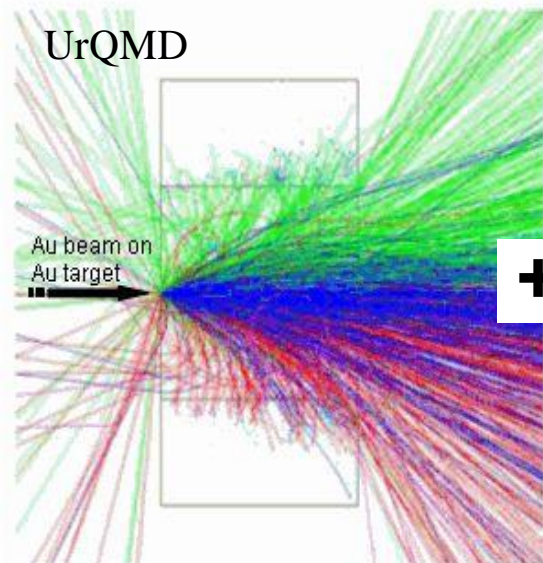
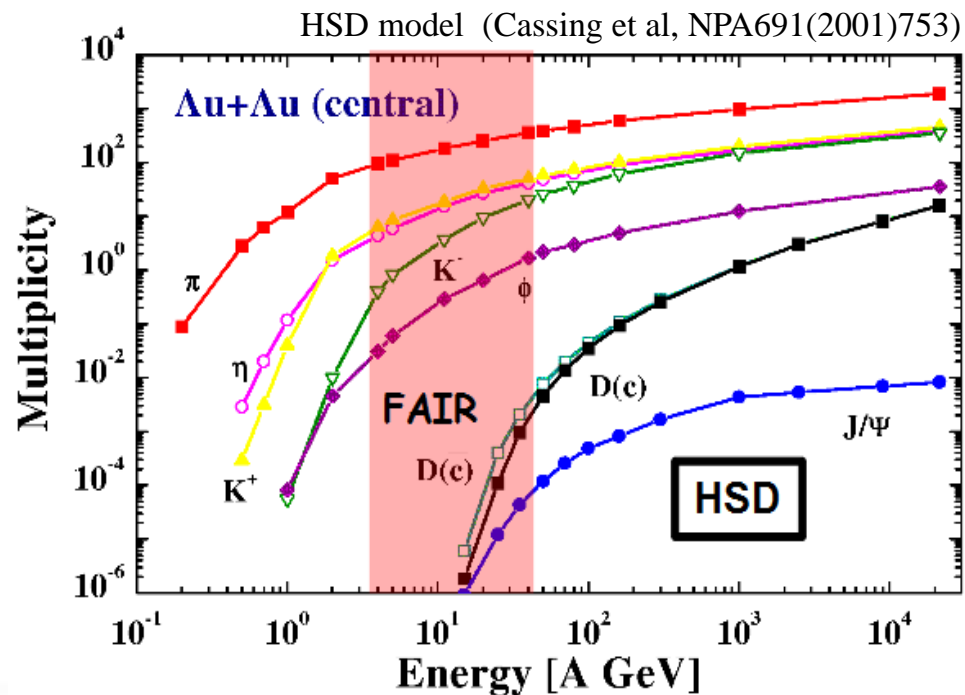


Open charm measurement with CBM

FAIR energies are at the kinematical threshold of open charm production:
 $M_{D^+} \sim 10^{-5}$ among ~ 1000 part./coll
 in central Au-Au coll.

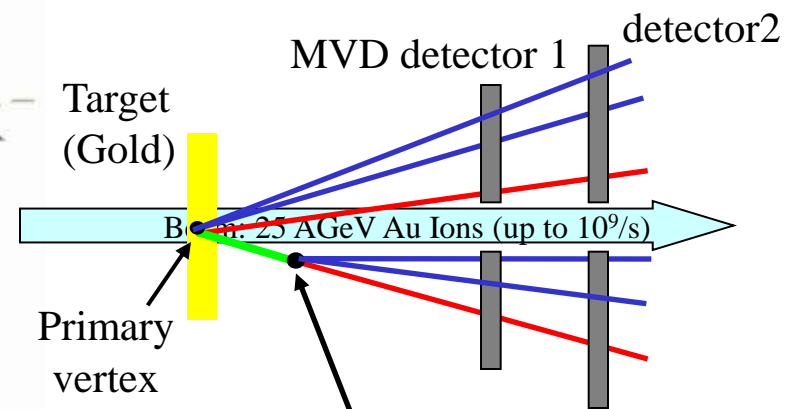
→

The Micro-Vertex Detector (MVD) will allow to disentangle open charmed and bulk particles

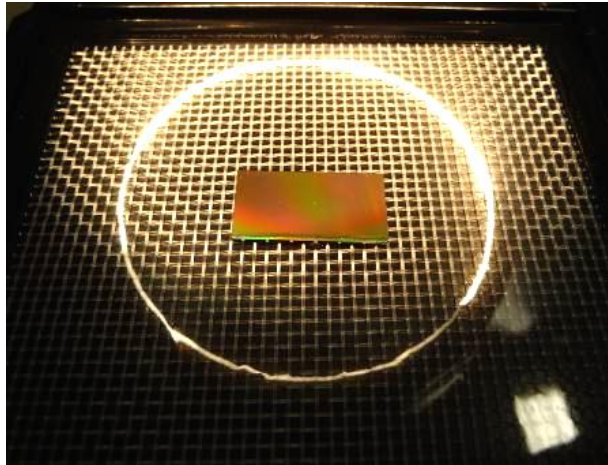


$$D^+ \rightarrow \pi^+\pi^+K^- \quad (c\tau = 317 \mu\text{m})$$

$$D^0 \rightarrow K^-\pi^+ \quad (c\tau = 124 \mu\text{m})$$



MAPS sensors for the MVD – R&D



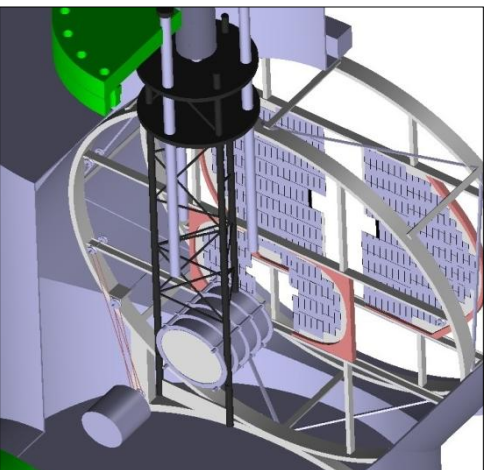
Monolithic Active Pixel Sensors
(MAPS, CMOS-based sensors)

- Invented by industry (digital camera)
- Modified for charged particle detection since 1999 by IPHC Strasbourg
- Also foreseen for ILC, STAR...

Rad. hard. and speed optimization on-going

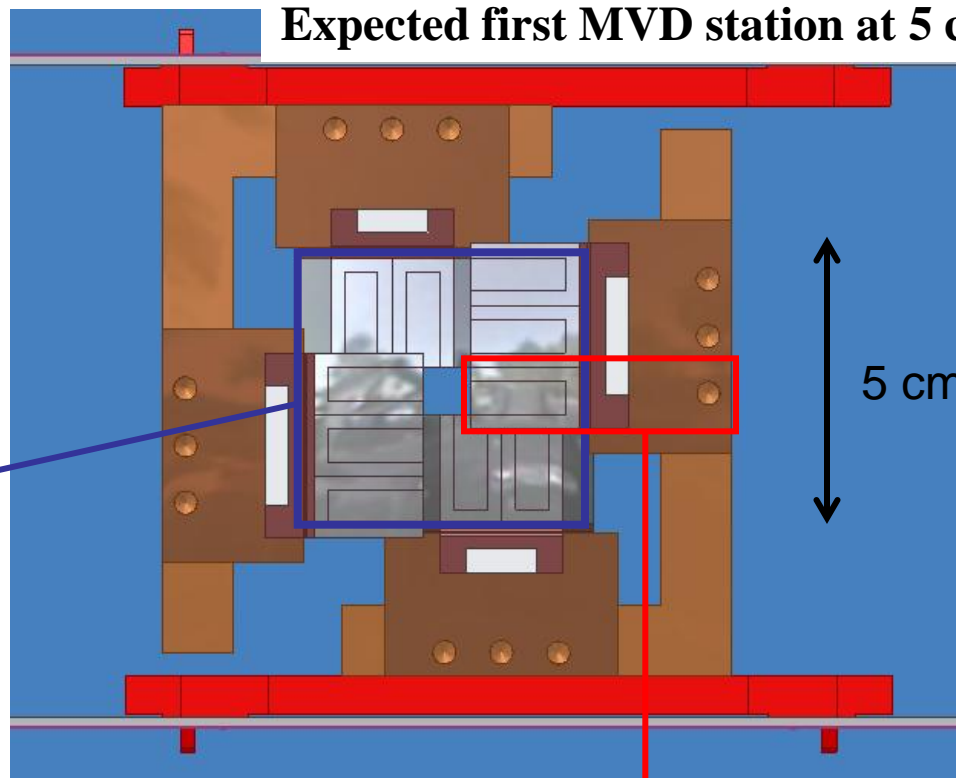
	CBM wish list	MAPS* (2003)	MAPS* (2009)	MIMOSA-26 Binary, \emptyset
Single point res.	$\sim 5 \mu\text{m}$	$1.5 \mu\text{m}$	$1 \mu\text{m}$	$4 \mu\text{m}$
Material budget	$< 0.3\% X_0$	$\sim 0.1\% X_0$	$\sim 0.05\% X_0$	$\sim 0.05\% X_0$
Rad. hard. non-io.	$> 10^{13} n_{\text{eq}}$	$10^{12} n_{\text{eq}}/\text{cm}^2$	$> 3 \times 10^{13} n_{\text{eq}}$	few $10^{12} n_{\text{eq}}$
Rad. hard. io	$> 3 \text{ Mrad}$	200 krad	$> 1 \text{ Mrad}$	$> 300 \text{ krad}$
Time resolution	$< 30 \mu\text{s}$	$\sim 1 \text{ ms}$	$\sim 25 \mu\text{s}$	$110 \mu\text{s}$

MAPS sensors for the MVD – integration



IKF - Frankfurt

Expected first MVD station at 5 cm



CBM-Acceptance
Low material

Bonds

Flex-Cable

Si

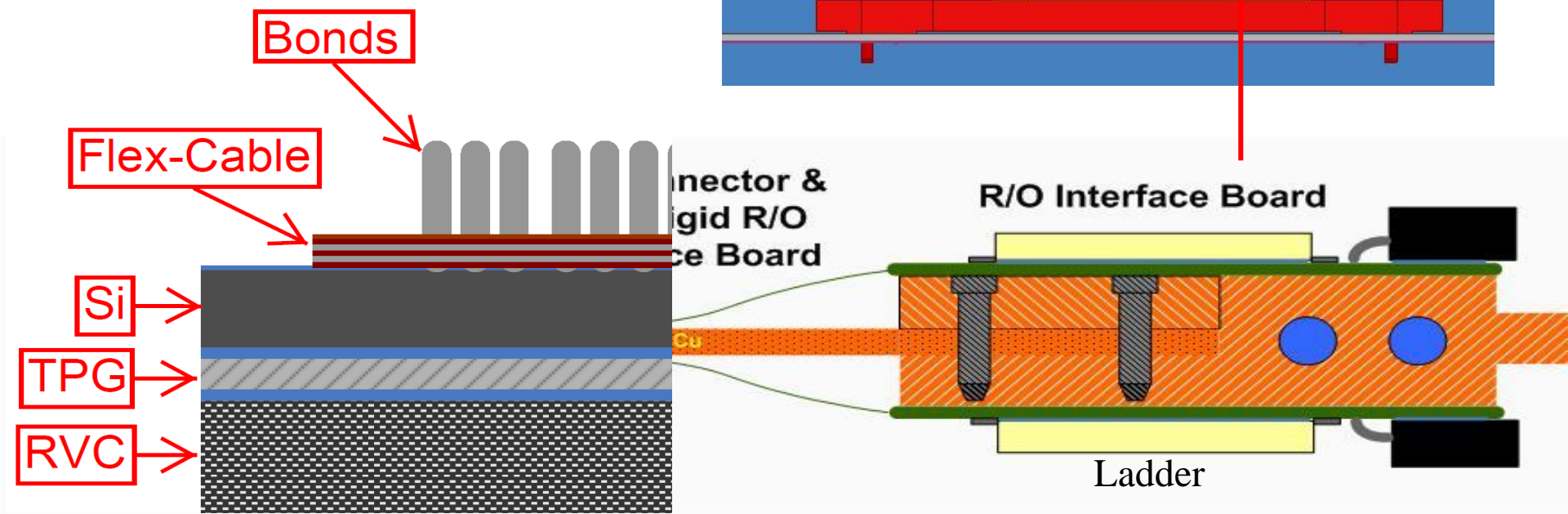
TPG

RVC

Connector &
Rigid R/O
Interface Board

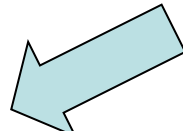
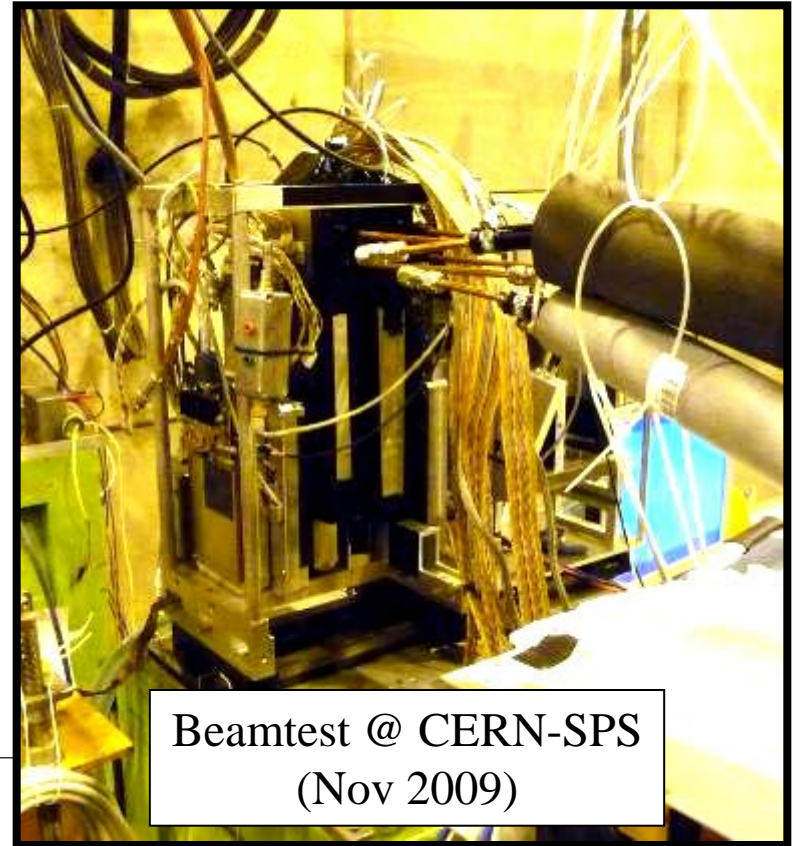
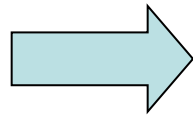
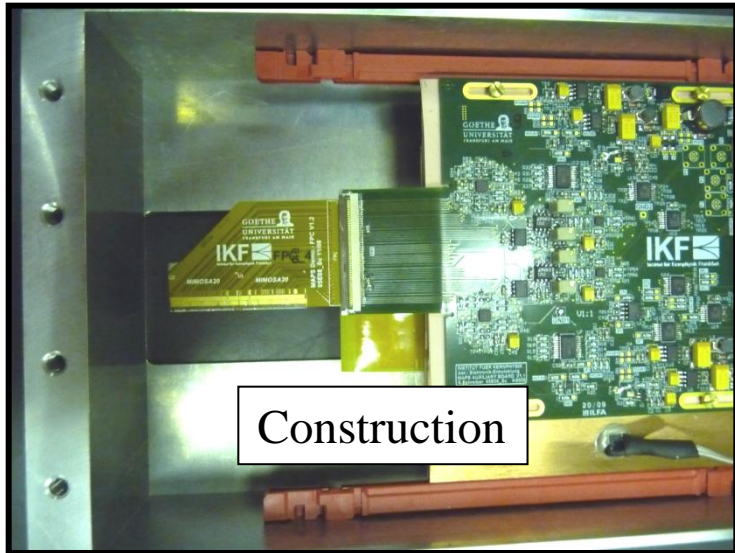
R/O Interface Board

Ladder

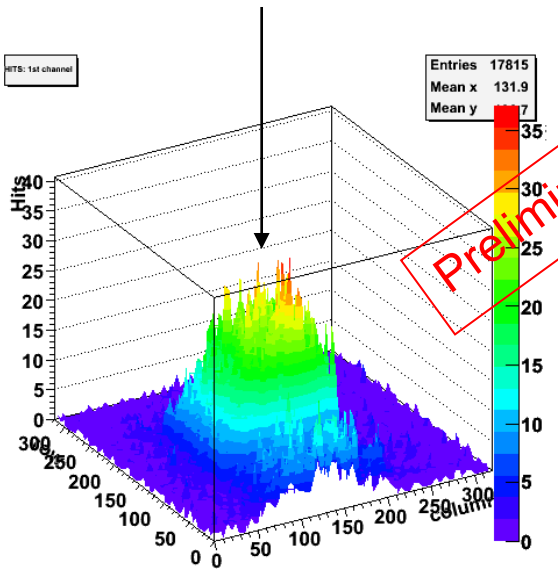


MAPS sensors for the MVD – integration (3)

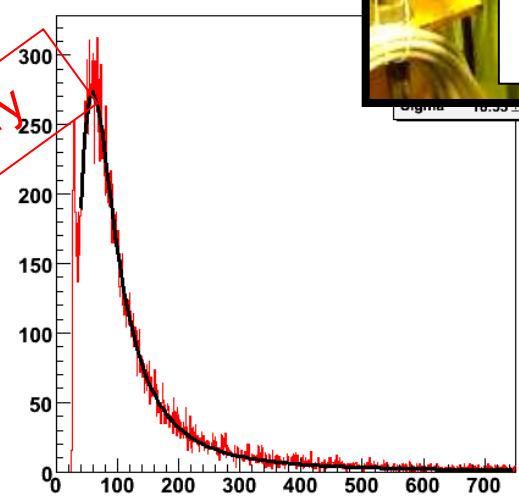
S. Amar-Youcef



Shadow of trigger scintillator



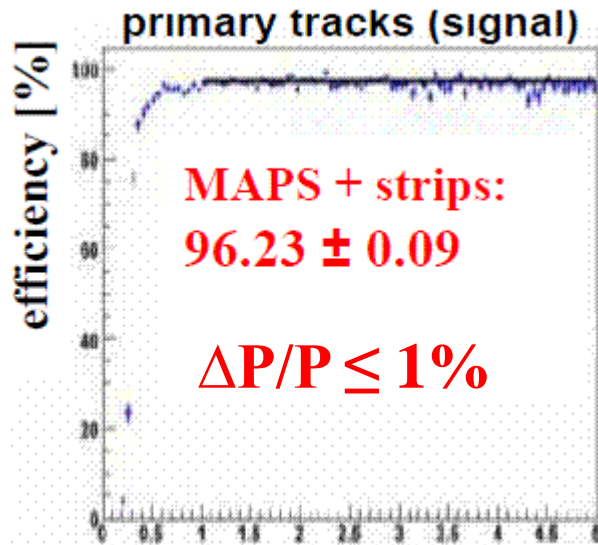
charge spectrum



Achievements:
System design validated
Good noise: 21 e⁻ ENC
Spatial res.: < 6 μm

Open charm measurement with CBM in the simulation

Tracking performances:



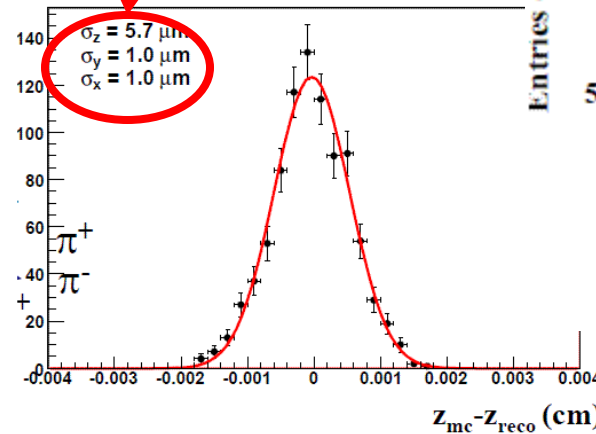
Topological cuts:

Track impact parameter

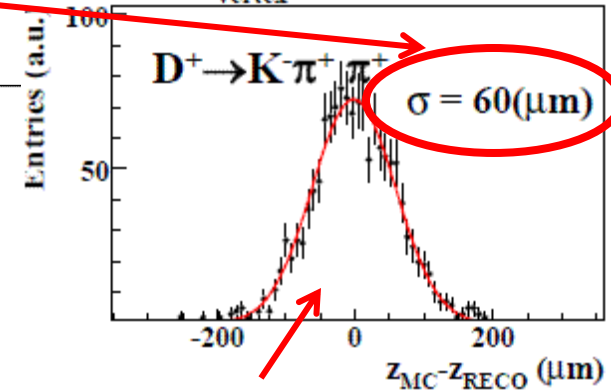
Vertex position, chi2



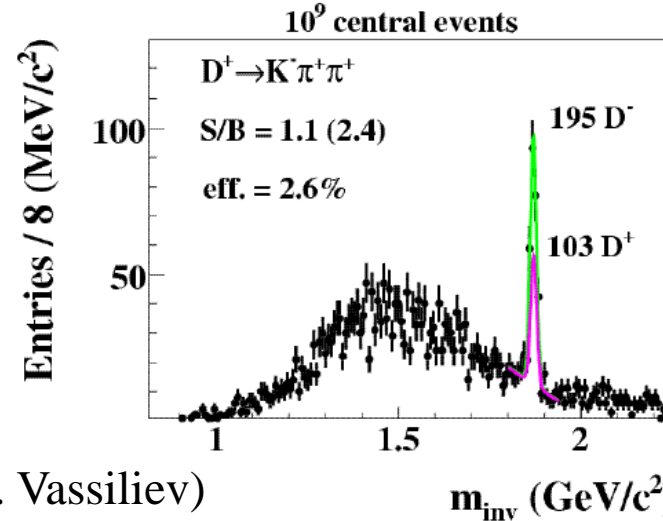
Primary & secondary vertex resolution



D^+ z_{vertex} resolution



with MAPS(150 μm),
ideal response, $\sigma_{sp} = 3\mu\text{m}$



(I. Vassiliev)

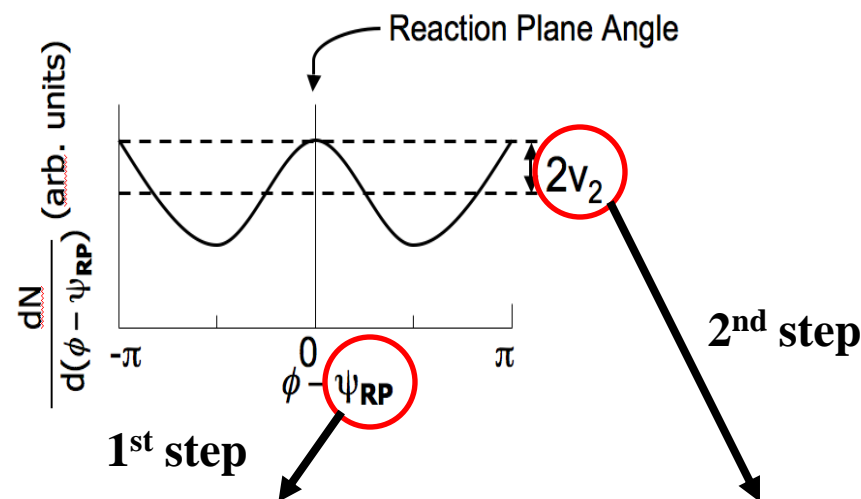
Central Au-Au @ 25 AGeV \rightarrow several 10^4 open charm particles per month (10^{11} central coll.)

Main steps of the v_2 measurement capability study with CBM

What is the expected precision for the measurement of open charm elliptic flow with CBM?

→ 2 limitations: **the reaction plane resolution**

the limited statistics of open charm particles



Evaluation of the reaction plane resolution:

- Use Au(25A GeV)-Au minimum bias events
→ transport through the appropriate detector(s)
- Systematic study of the reaction plane resolution
→ **b and Ebeam dependence**

Reconstruction of v_2 for the D^+ :

Assume the reaction plane resolution evaluated in the 1st step

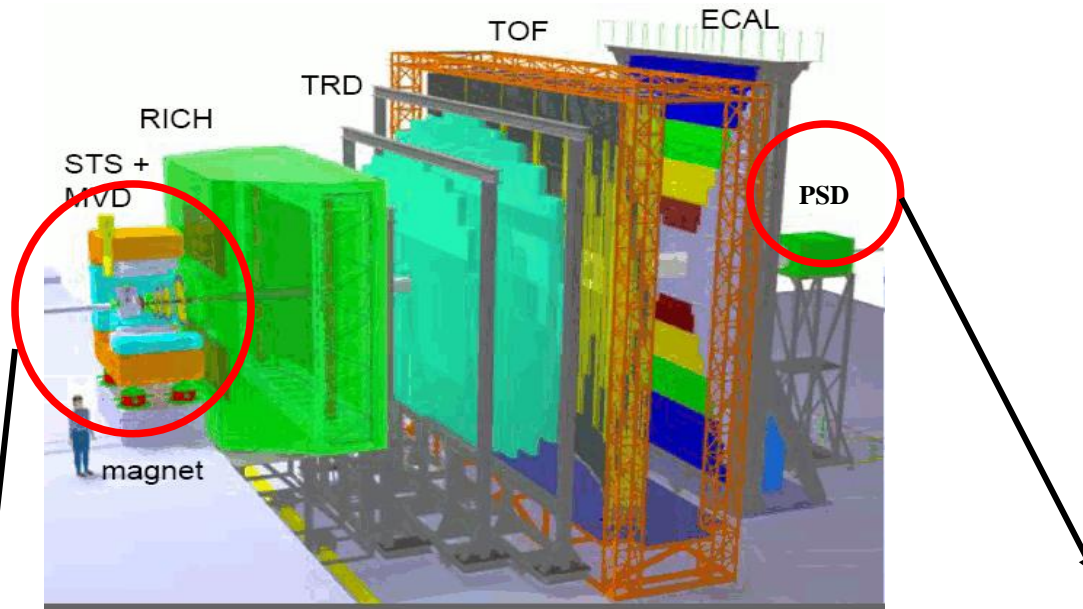
this will give a estimate of the v_2 statistical error!

$\sigma_{\Phi_{RP}}$

$M_D \sim 10^{-5}$ for 25 A GeV Au-Au central coll.

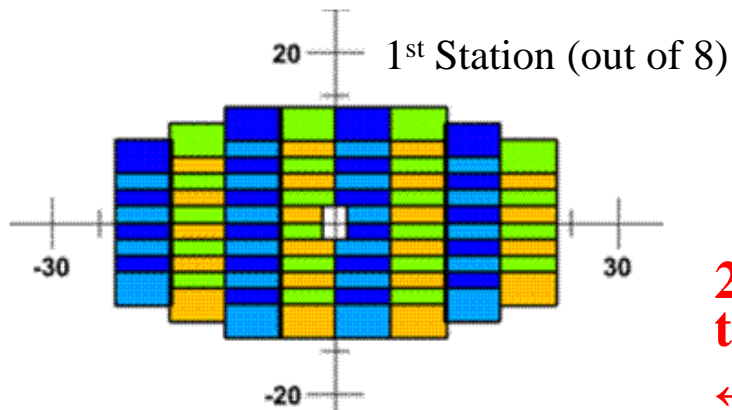
Reconstruction of the reaction plane

- Experimental set-up -



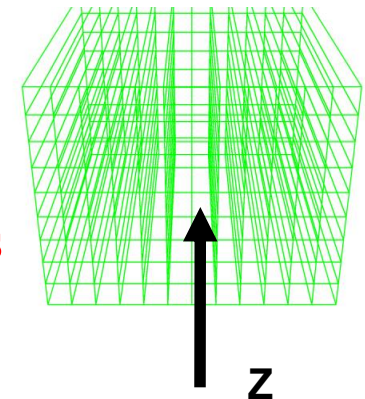
Silicon Tracking System (STS)

Projectile Spectator Detector (PSD)
CBM-PSD-note-2006-001



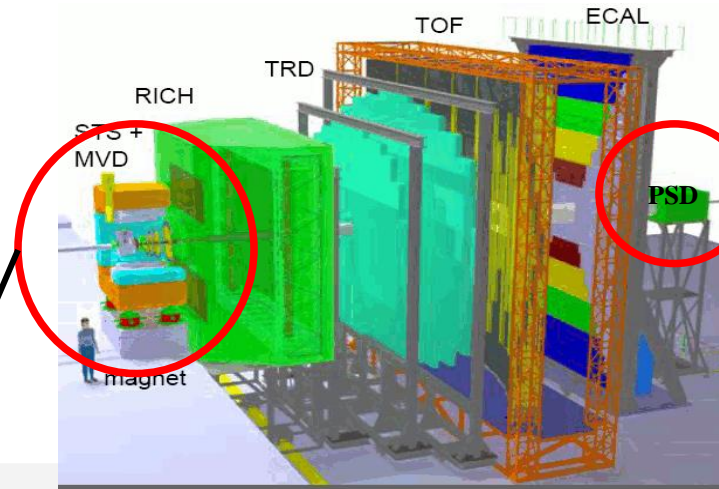
Heuser, ULISI, Feb 2010

**2 independent sub-detectors
to measure the event plane
↔ cross-check of several
estimates of v_2 !**

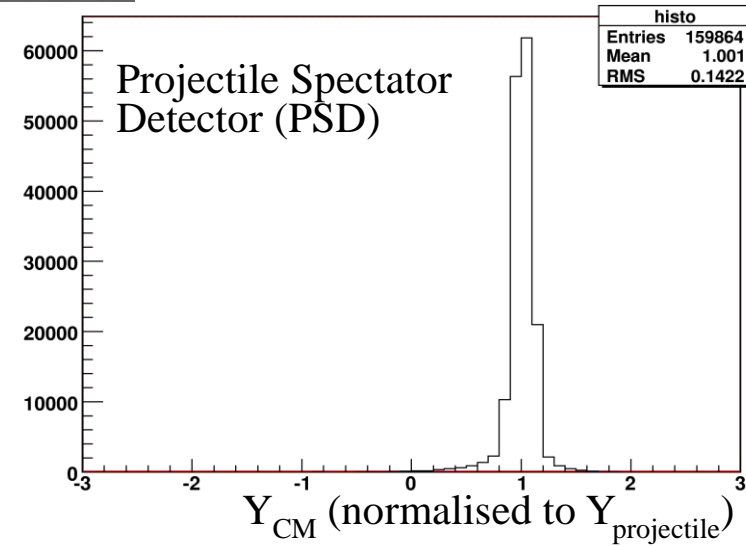
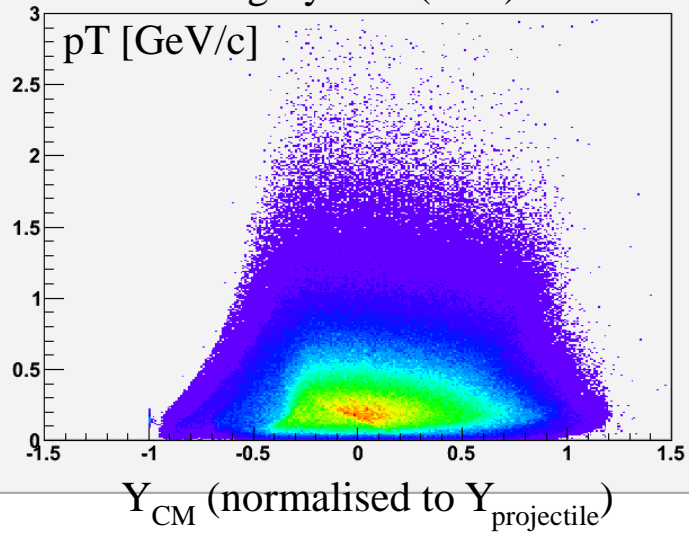


Reconstruction of the reaction plane

- Detector acceptance -



Silicon Tracking System (STS)

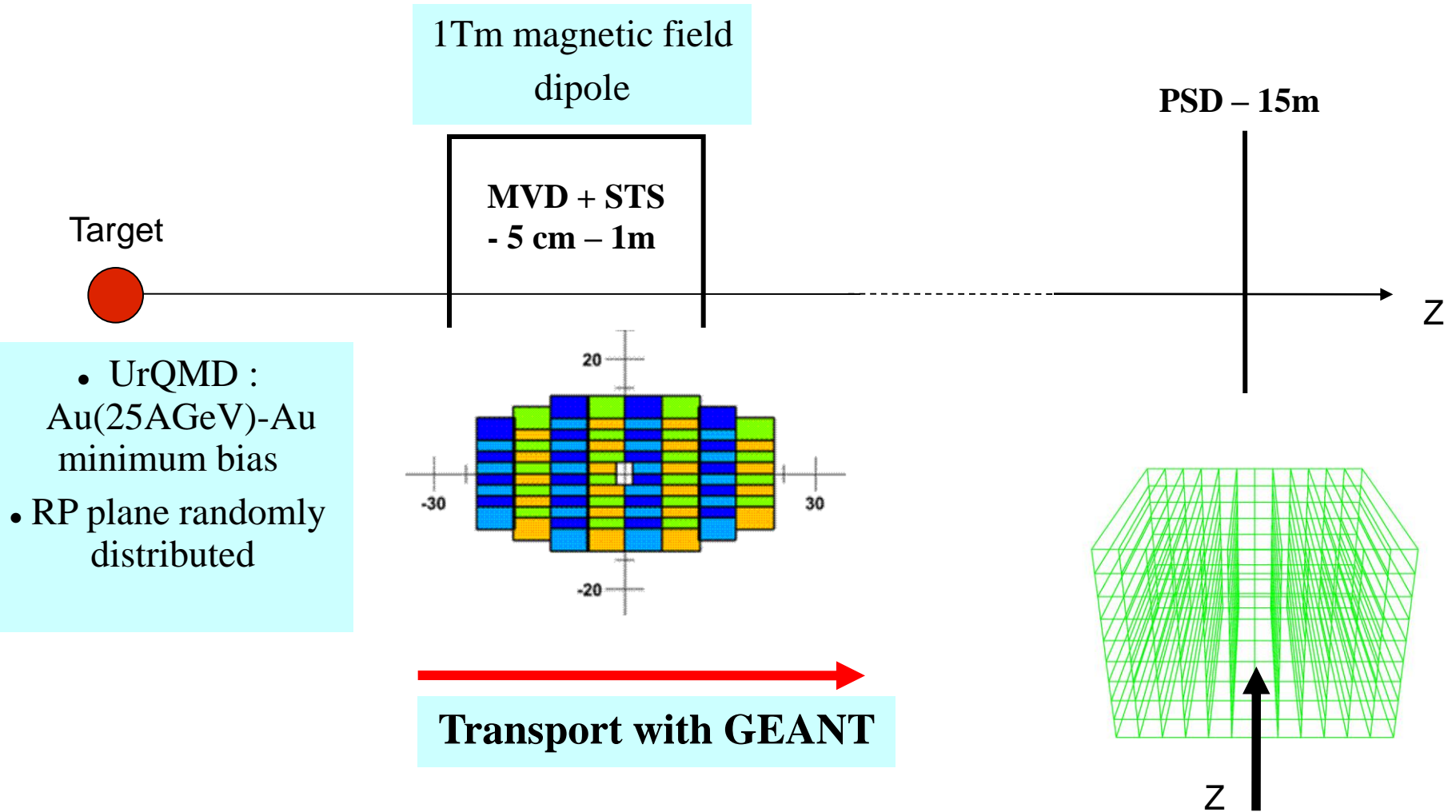


STS → mid-Y / forward-Y coverage → n = 2 and 1

PSD → forward-Y coverage → n = 1

Reconstruction of the reaction plane

- Simulation set-up -



Reconstruction of the reaction plane

- Method for the STS -

Poskanzer and S. Voloshin,
arXiv:nucl-ex/9805001

Flow vector Q:

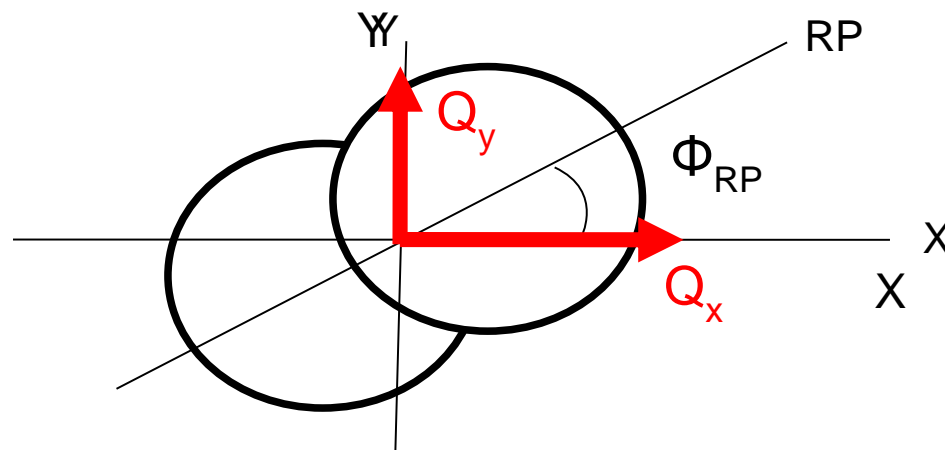
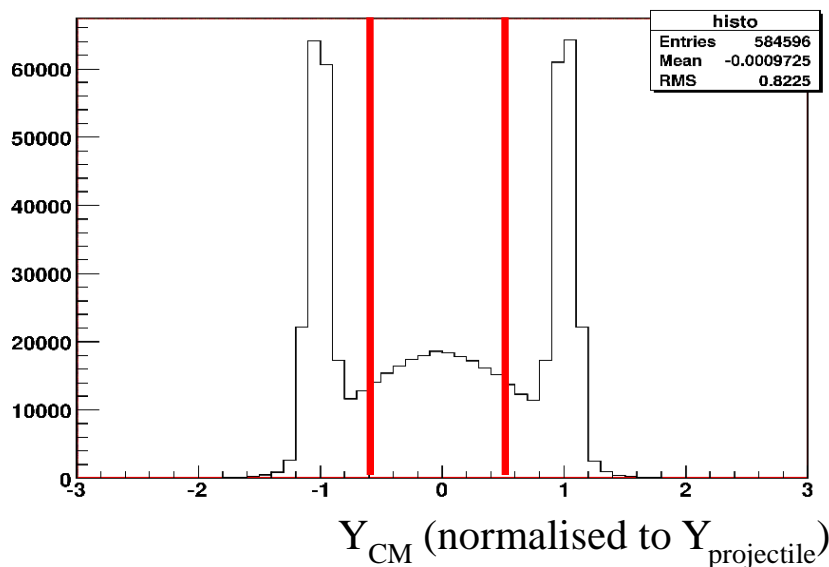
$$Q_x = \sum_i w^i \cdot p_T^i \cdot \cos(n \cdot \Phi^i)$$

$$Q_y = \sum_i w^i \cdot p_T^i \cdot \sin(n \cdot \Phi^i)$$

$$\Phi_{RP}^{reco} = 1/n \cdot \tan^{-1} (Q_y / Q_x)$$

p_T^i : transverse momentum of part. i

Φ^i : azimuth of the part. i



Also: anti-flow of pions

$$\text{For } n = 1: \Phi_{\text{pion}} \rightarrow \Phi_{\text{pion}} + \pi$$

For n = 1: W = -1 0 1

For n = 2: W = 1

Reconstruction of the reaction plane

- Method for the PSD -

Flow vector Q :

Gravity center of the energy deposited in the PSD modules:

$$Q_x = \sum_i R_{\text{module}}^i \cdot E_{\text{module}}^i \cdot \cos(\Phi_{\text{module}}^i)$$

$$Q_y = \sum_i R_{\text{module}}^i \cdot E_{\text{module}}^i \cdot \sin(\Phi_{\text{module}}^i)$$

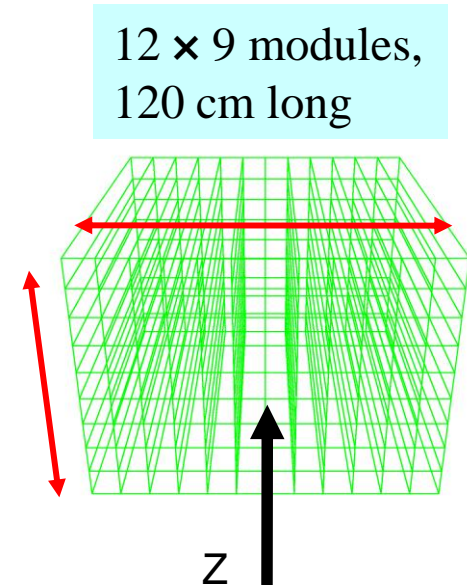
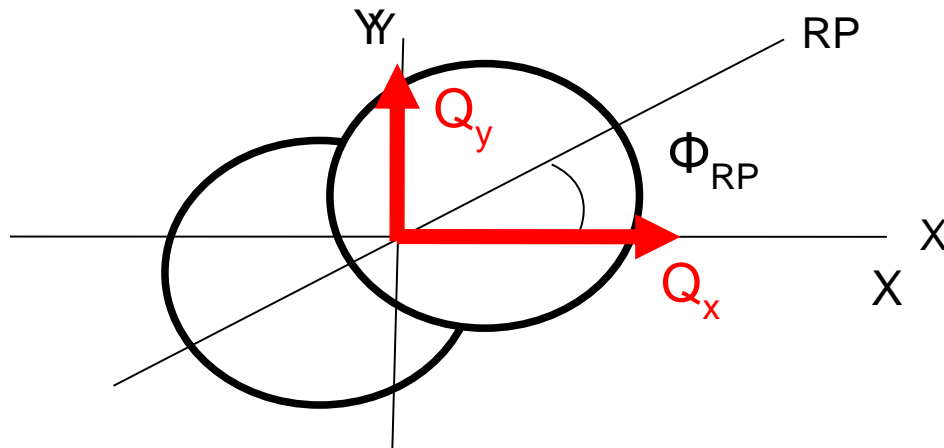
$$\Phi_{\text{RP}}^{\text{reco}} = \tan^{-1}(Q_y / Q_x)$$

Poskanzer and S. Voloshin,
arXiv:nucl-ex/9805001

E_{module}^i : deposited energy in module I

R_{module}^i : radius of module I \leftrightarrow distance from (0, 0)

Φ_{module}^i : azimuth of the module i



Reconstruction of the reaction plane

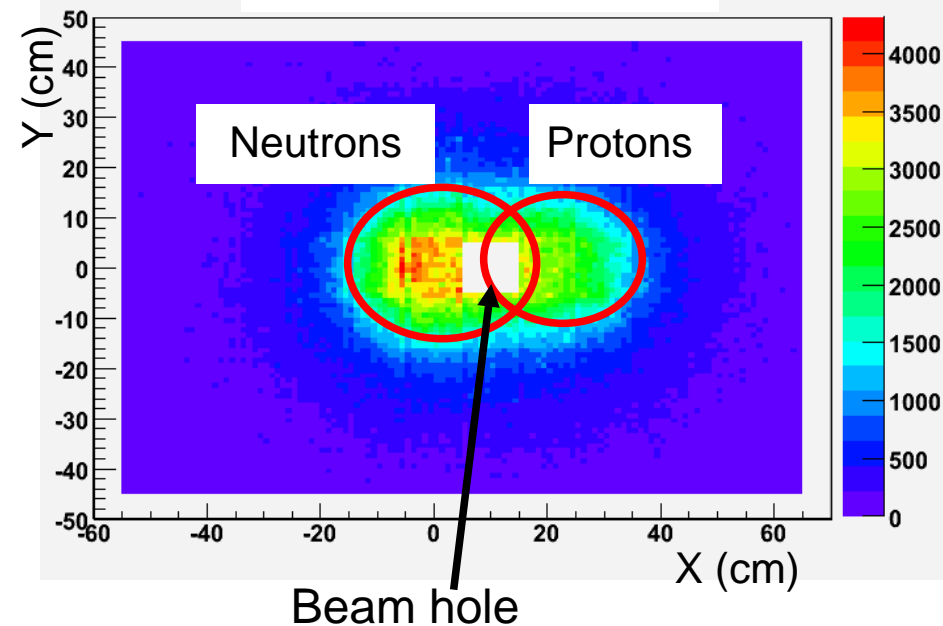
- PSD azimuthal asymmetry and shift of the protons-

- **T**he PSD is shifted along $X > 0$ to place the **beam hole** at the beam spot position
- **T**he spectator protons are deflected by the magnetic dipole field
- **W**e used one of the simplest technique to **flatten the reconstructed event plane distribution**:

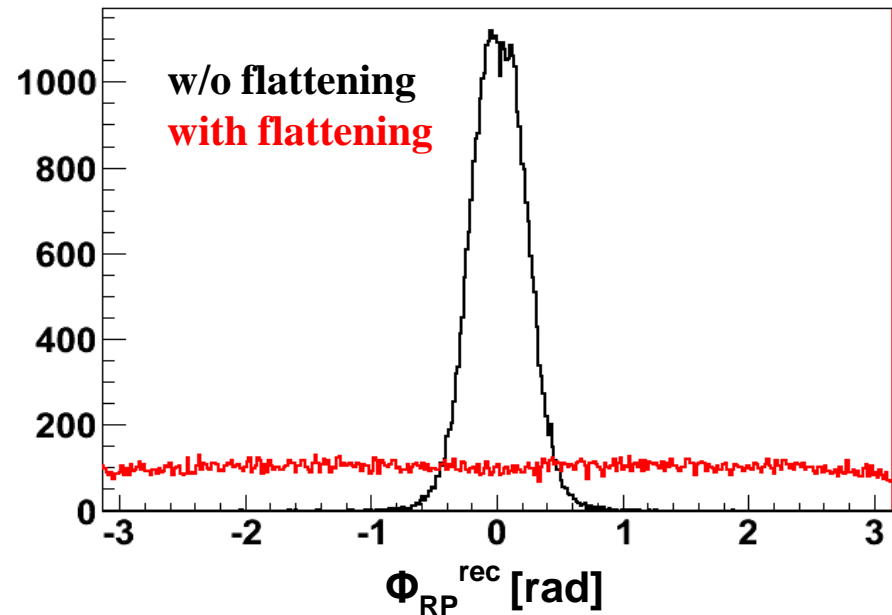
$$Q_x = \sum_i E_{\text{module}} \cdot X_{\text{module}} - \langle \sum_i E_{\text{module}} \cdot X_{\text{module}} \rangle$$

- **T**he phi weight method has been used but not conclusive results ...
- **O**ther flattening methods: shifting method, etc
(A. Poskanzer and S. Voloshin, arXiv:nucl-ex/9805001, J. Barrette et al. arXiv:nucl-ex/9707002)

Hit distribution in the PSD



Reconstructed event plane



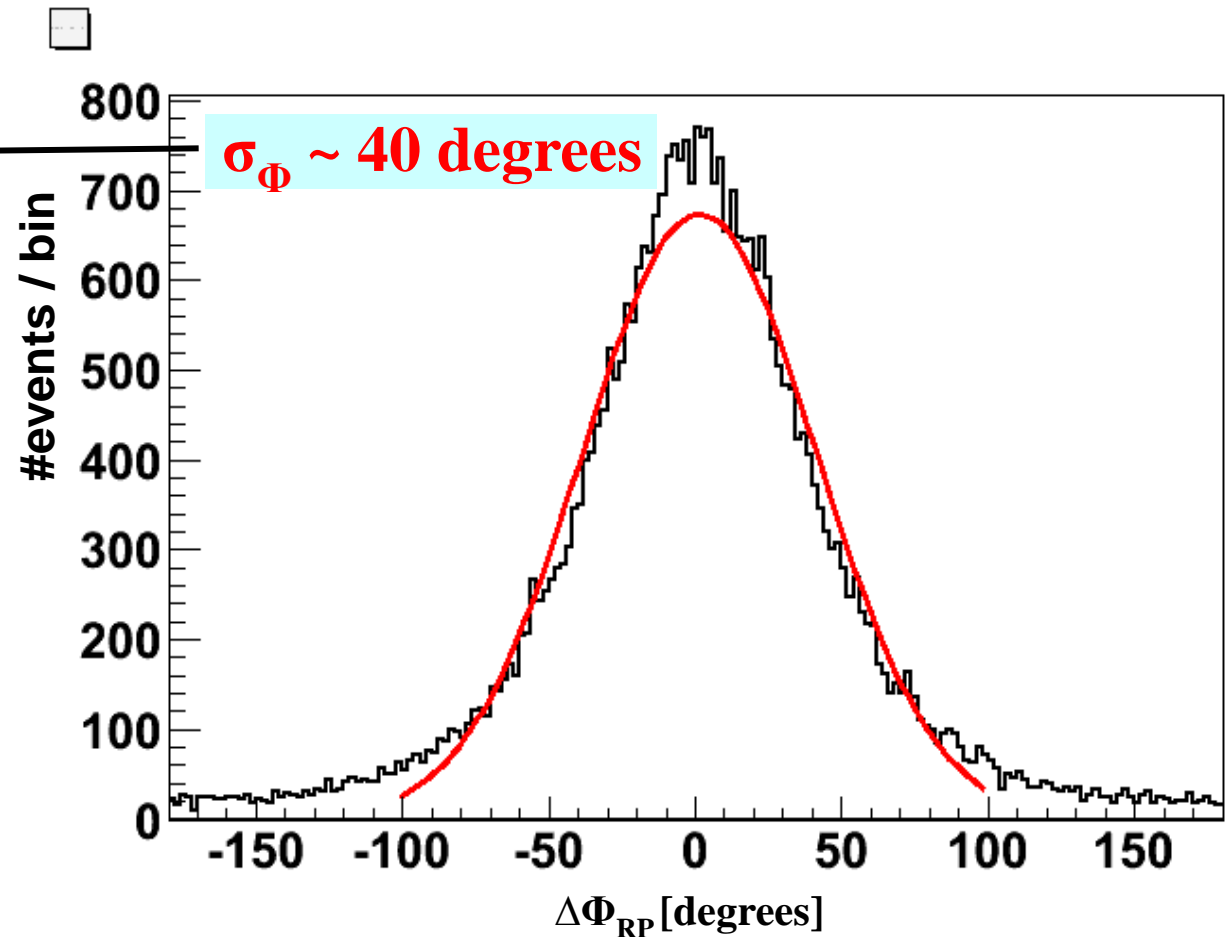
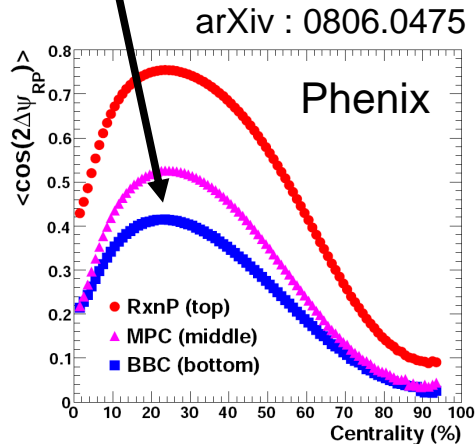
Reconstruction of the reaction plane

- $\Delta\Phi_{\text{RP}} = \Phi_{\text{RP}}^{\text{reco}} - \Phi_{\text{RP}}^{\text{true}}$ distribution for the PSD event plane -

- Au(25A GeV)-Au with b between 6 and 9 fm
- $\Delta\Phi_{\text{RP}} \sim \text{Gaussian}(0, \sigma_{\Phi})$

$\langle \cos(2 \Delta\Phi_{\text{RP}}) \rangle_{\text{event}} \sim 0.3$

For comparison:



Reconstruction of the reaction plane

Ideal calculation of v_2 :

$$v_2 = \frac{\langle p_x^2 \rangle - \langle p_y^2 \rangle}{\langle p_x^2 \rangle + \langle p_y^2 \rangle} = \langle \cos(2 \times (\phi - \Psi_{RP})) \rangle$$

True reaction plane

Real calculation of v_2 :

$$v_{2obs} = \langle \cos(2 \times (\phi - \Psi_n)) \rangle$$

Measured event plane (of order n)

Correction:

$$\langle \cos(2 \times (\phi - \Psi_n)) \rangle = \langle \cos(2 \times (\phi - \Psi_{RP})) \rangle \langle \cos(2 \times (\Psi_{RP} - \Psi_n)) \rangle$$

$$\longrightarrow \underbrace{\langle \cos(2 \times (\phi - \Psi_{RP})) \rangle}_{\text{True } v_2} = \underbrace{\langle \cos(2 \times (\phi - \Psi_n)) \rangle}_{v_{2obs}} / \boxed{\langle \cos(2 \times (\Psi_{RP} - \Psi_n)) \rangle}$$

True v_2

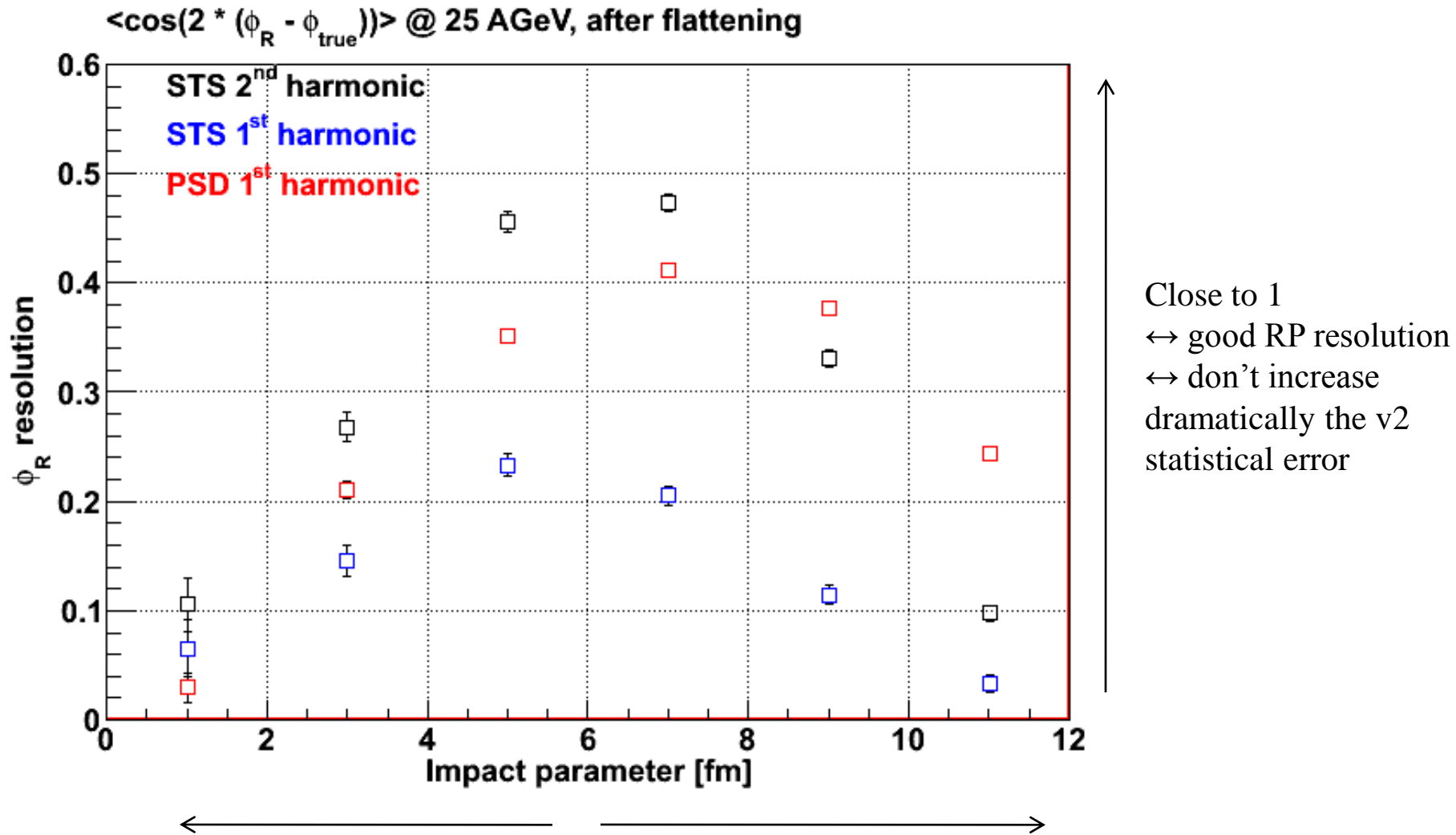
v_{2obs}

Now on, reaction plane resolutions expressed in terms of $\langle \cos(2 \times (\Psi_{RP} - \Psi_n)) \rangle$

Side remark: the statistical error on the true v_2 is proportional to $1 / \langle \cos(2 \times (\Psi_{RP} - \Psi_n)) \rangle$

Reconstruction of the reaction plane

- Impact parameter (b) dependence of the event plane resolution -



← Toward central coll.

→ low initial eccentricity

↔ low anisotropy in P devpt

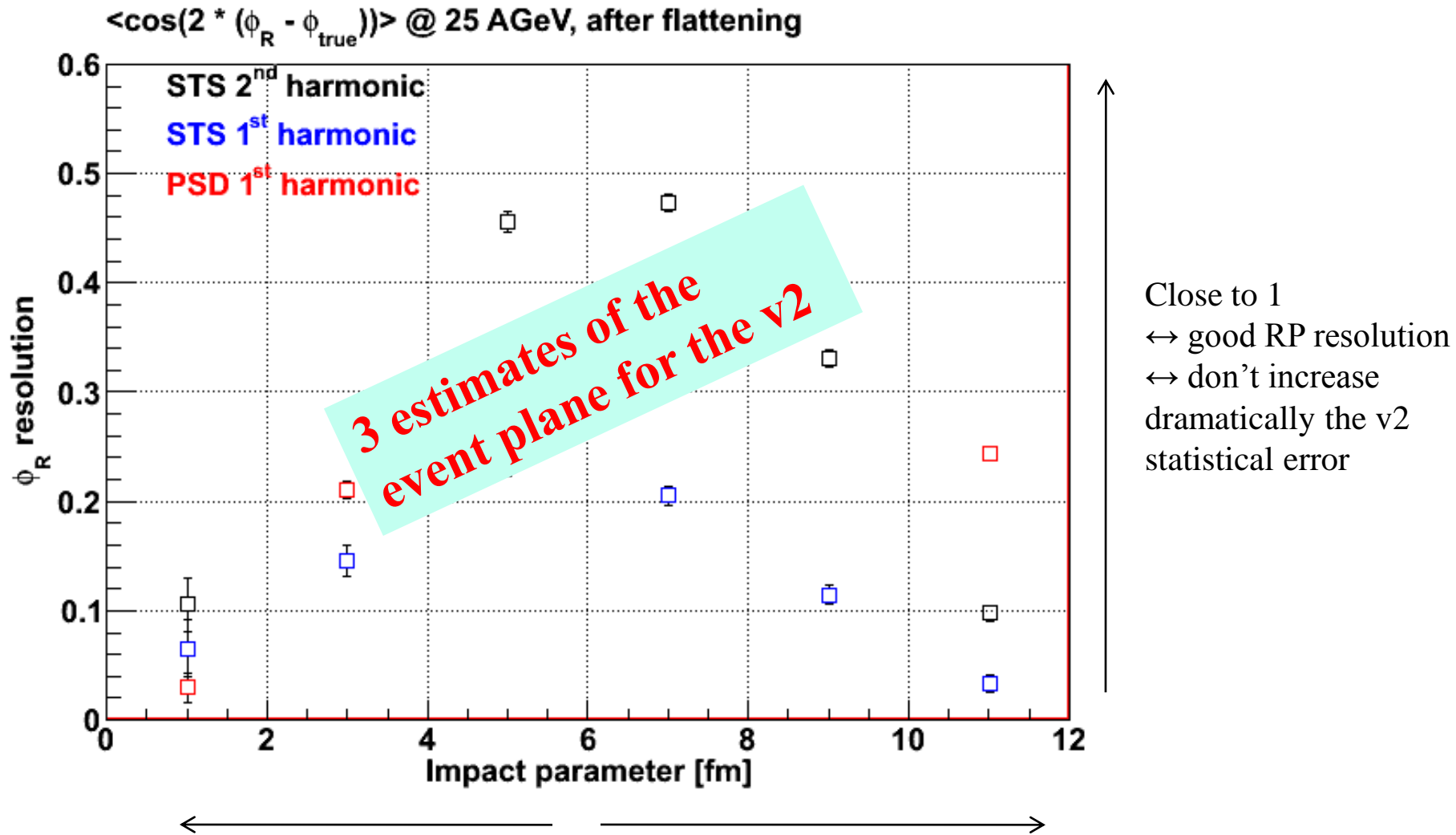
→ Toward peripheral coll.

→ low energy density (low P devpt)

→ for STS: low acceptance

Reconstruction of the reaction plane

- Impact parameter (b) dependence of the event plane resolution -

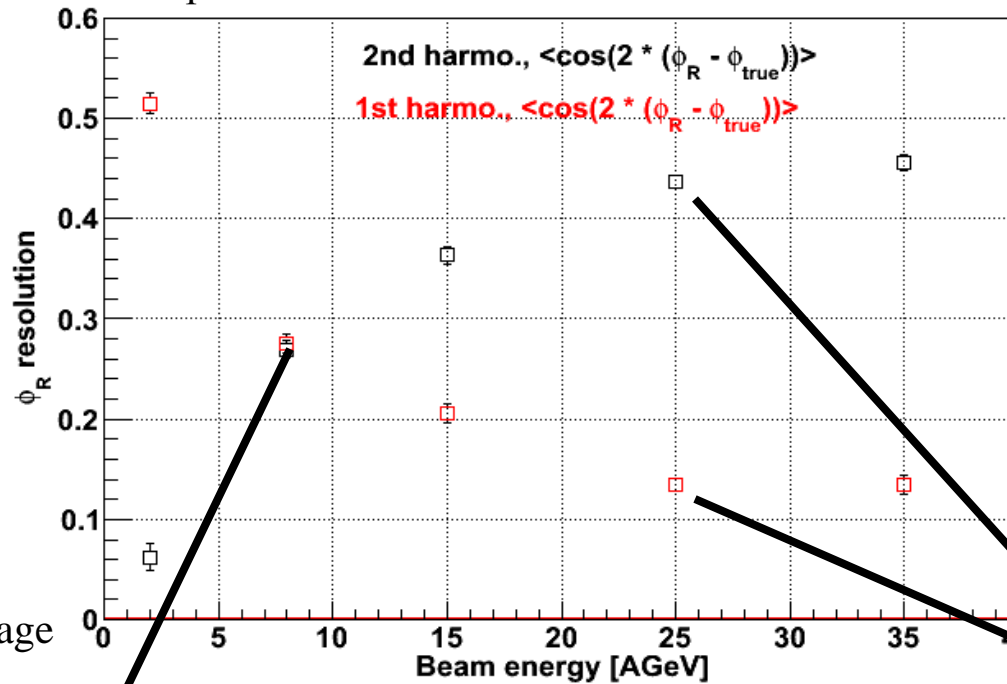


← Toward central coll. → low initial eccentricity
 ↔ low anisotropy in P devpt
 → Toward peripheral coll. → low energy density (low P devpt)
 → for STS: low acceptance

Reconstruction of the reaction plane

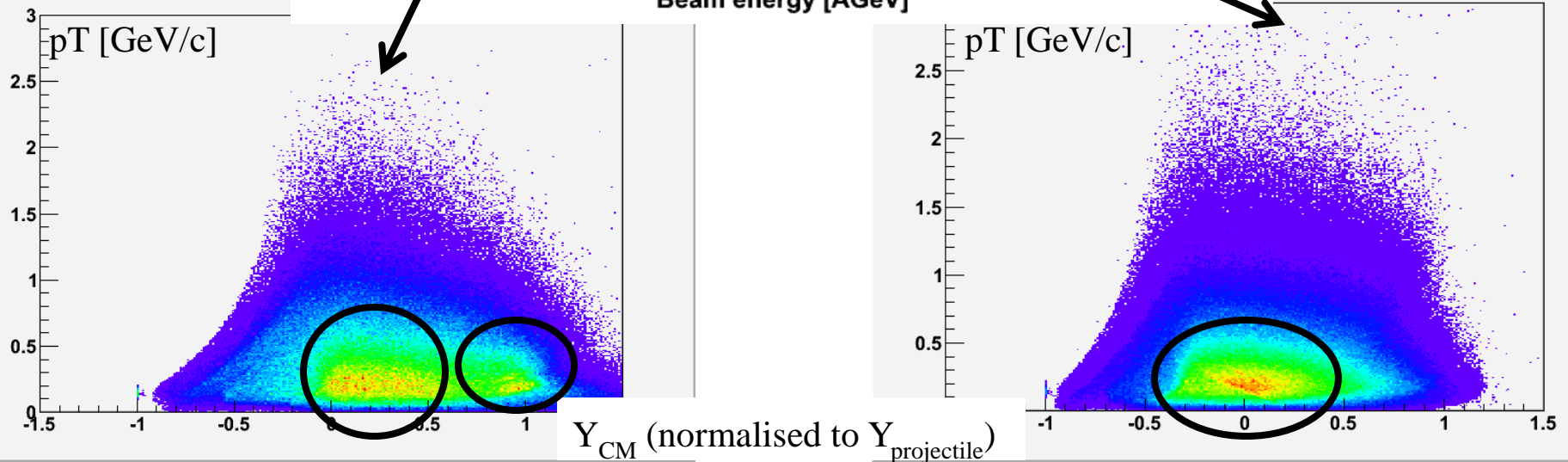
- Beam energy (E_{beam}) dependence for the STS event plane resolution -

Acceptance \leftrightarrow 4 hits in STS



For $E \leq 15$ A.GeV:
Good FMR coverage
And partial FR coverage

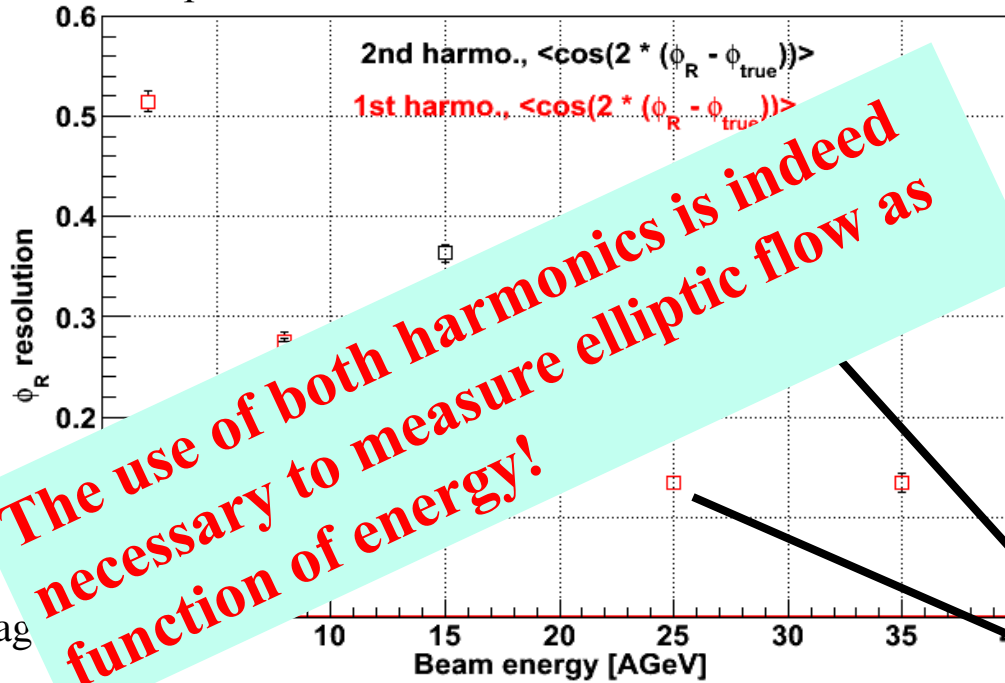
For $E > 15$ A.GeV:
Good MR coverage
But bad FR coverage ...



Reconstruction of the reaction plane

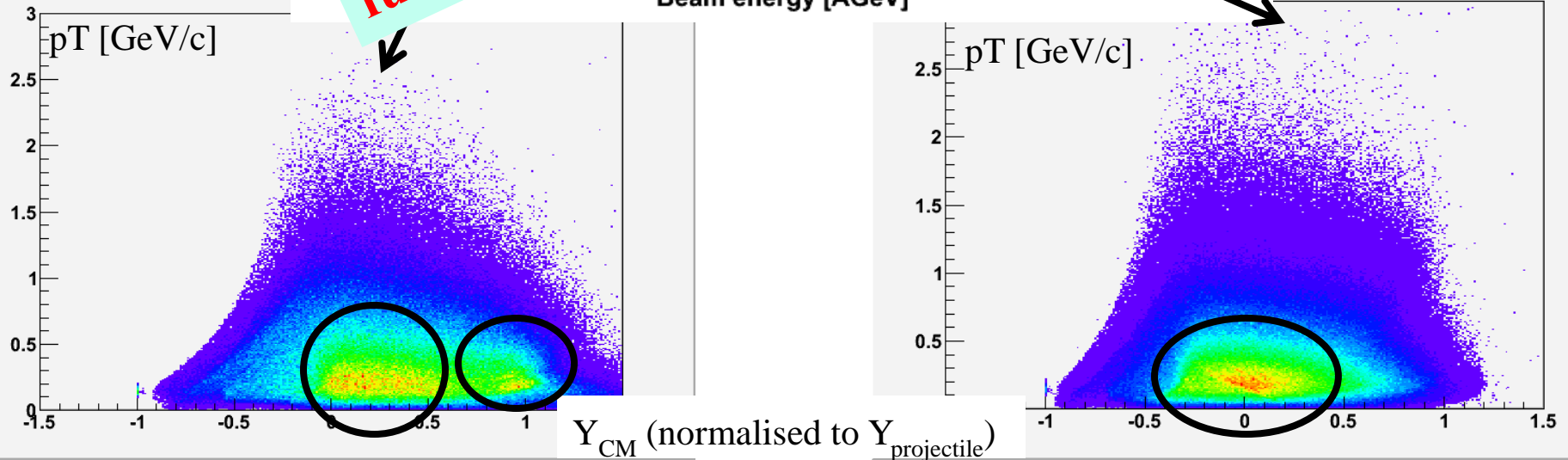
- Beam energy (E_{beam}) dependence for the STS event plane resolution -

Acceptance \leftrightarrow 4 hits in STS



For $E > 15$ A.GeV:
Good MR coverage
But bad FR coverage ...

For $E \leq 15$ A.GeV:
Good FMR coverage
And partial FR coverage



Reconstruction of the reaction plane

- Experimental evaluation – sub-event method -

- Create 2 sub-events with equal multiplicity
- Calculate the event plane for each of them
- Flatten the 2 resulting event planes ...
- And use the correlation between them

$$\langle \cos[n(\Psi_n^a - \Psi_n^b)] \rangle$$

- The resolution of each sub-event plane is:

$$\text{resSub} \equiv \langle \cos(n(\Psi_n^a - \Psi_{RP})) \rangle = \sqrt{\langle \cos(n(\Psi_n^a - \Psi_n^b)) \rangle}$$

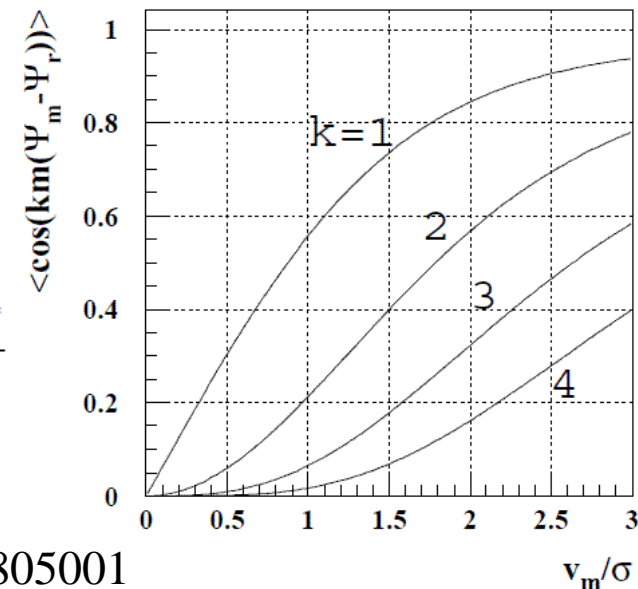
- This can be used in 1st approx. to evaluate the resolution of the full event plane (used to determine the v_2 parameter):

$$\text{res} \equiv \langle \cos(n(\Psi_n - \Psi_{RP})) \rangle \leq \sqrt{2} \text{resSub}$$

- Other method:

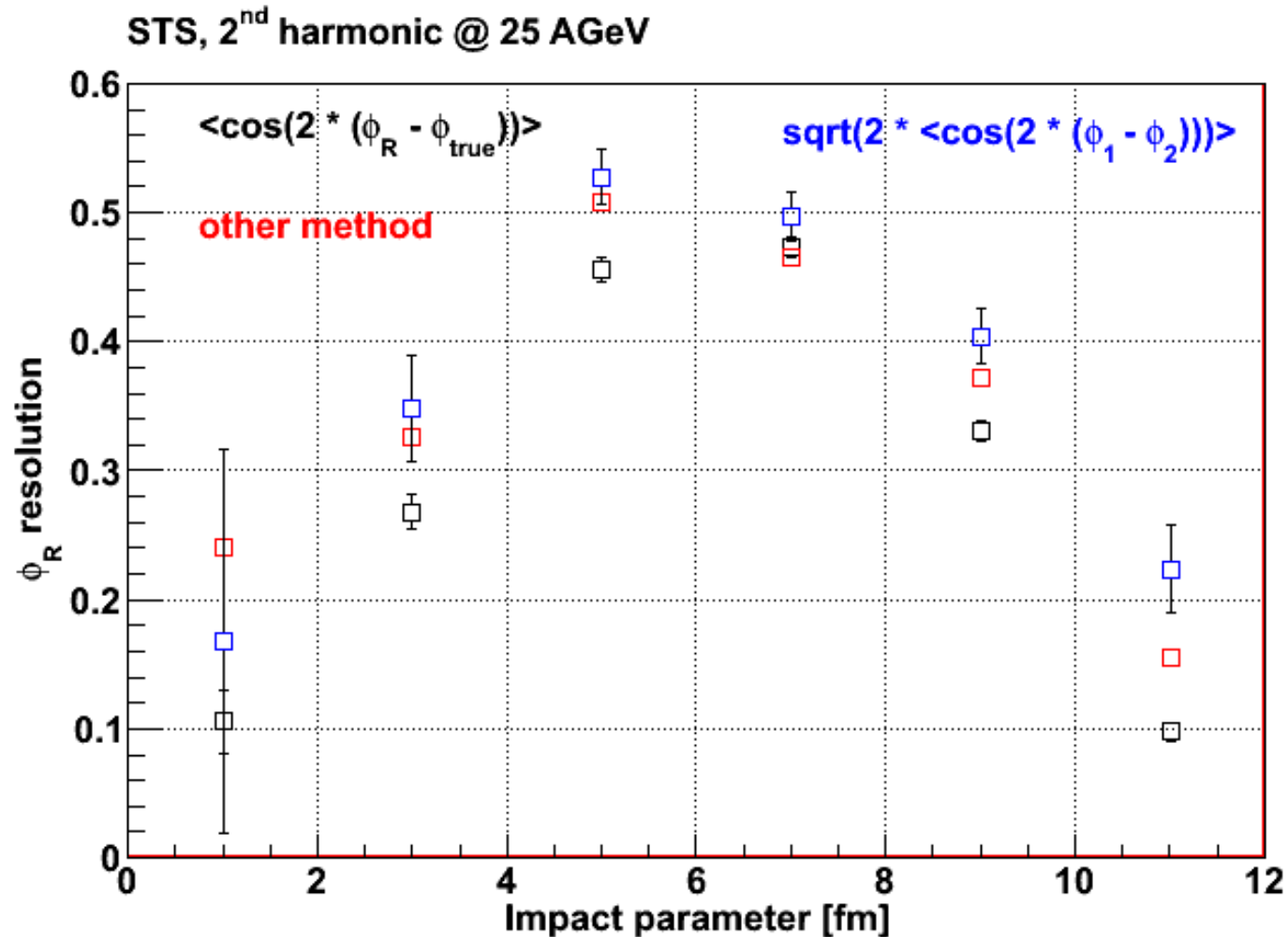
$$\frac{N_{events}(m|\Psi_m^a - \Psi_m^b| > \pi/2)}{N_{total}} = \frac{e^{-\chi_m^2/4}}{2}$$

$$\chi_m = v_m \sqrt{2N}$$



Reconstruction of the reaction plane

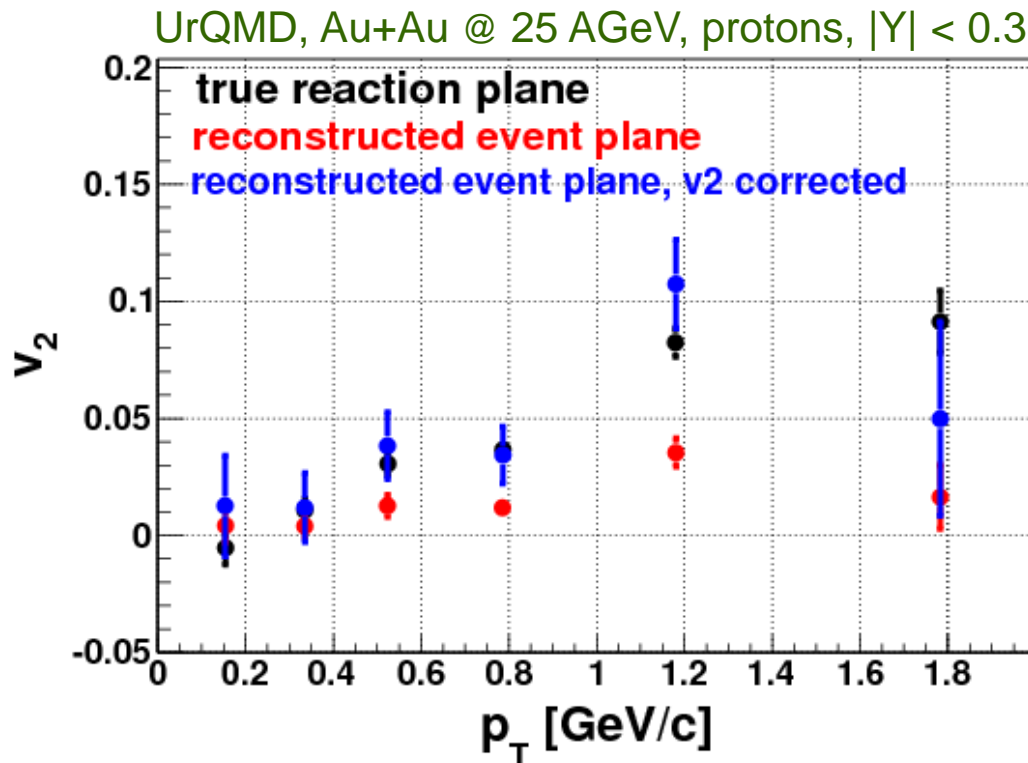
- Experimental evaluation – sub-event method -



- The experimentally determined event plane resolution should coincide with the actual event plane resolution reasonably well

Elliptic flow of bulk particles (protons, pions)

- For each event with b in $[6, 9]$ fm:
 - the event plane is evaluated with the PSD
 - v_2 of bulk particles is calculated with the MVD+STS: here the protons at mid-rapidity ($|Y| < 0.3$)
- The azimuthal distribution of the particles is taken relative to the reconstructed event plane
- $v_2 = \langle \cos\{2(\Phi - \Phi_{\text{RP}}^{\text{reco}})\} \rangle$
- $v_2^{\text{corr}} = v_2 / \text{Corr}$ $\text{Corr} = \langle \cos\{2(\Phi_{\text{RP}}^{\text{reco}} - \Phi_{\text{RP}}^{\text{true}})\} \rangle \sim 0.3$
- After a correction is applied, the v_2 reconstructed values fit the true values within statistical errors

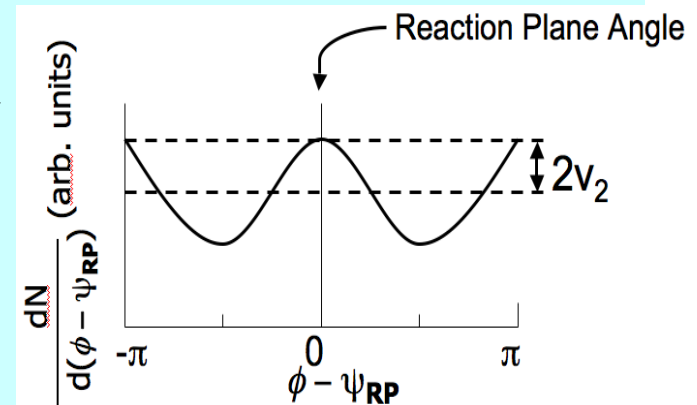


Only 20K minimum bias events
in the simulation!

v_2 reconstruction for the D^+

- Simulation method -

- **Generate D^+ mesons at mid-rapidity**
- **Assume this sample to be reconstructed D^+**
 - **very large coverage at mid-rapidity with CBM (confirmed by simulations)**
 - **very low background contamination**
- **Thermal distribution: $P(p_T) = p_T \cdot e^{-\sqrt{m^2/T}}$, $T = 200$ MeV, $m_{D^+} = 1.87$ GeV/c²**
- **Assume p_T linear dependence of v_2 : $v_2 = 0.03 \times p_T$, $v_2 = 0.05 \times p_T$, etc**
- **Azimuthal anisotropy**
 - $\Phi - \Phi_{RP}$ distributed according to the function $dN/d\Phi$
- **Reconstructed event plane uncertainty**
 - $\Delta\Phi_{RP} \sim \text{Gaus}(0, \sigma_{RP})$, $\sigma_{RP} = 40$ degrees

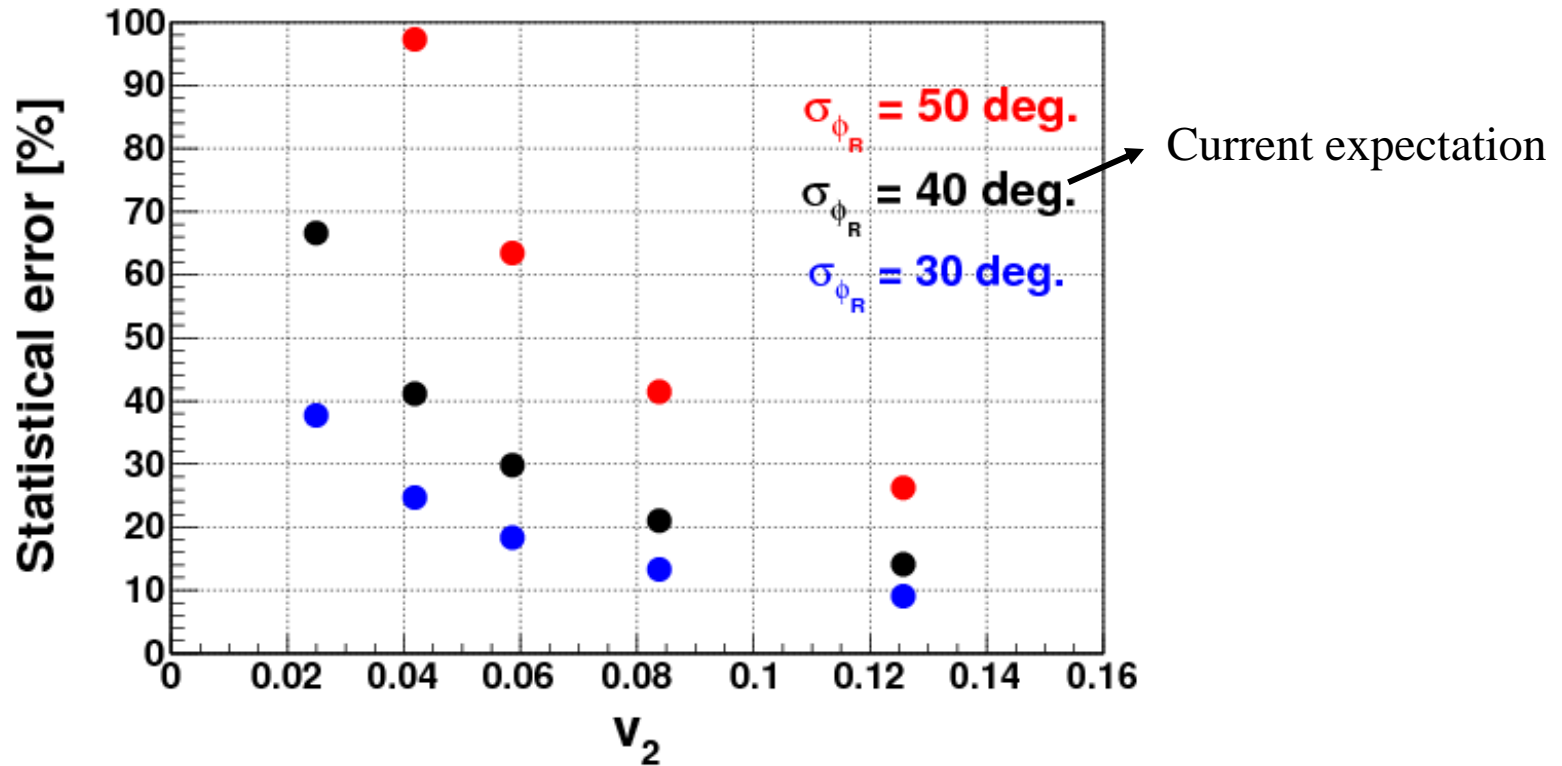


Current expectation

v_2 reconstruction for D^+

- Statistical errors on integrated v_2 -

We consider a sample of 10K reconstructed D^+ meson (~ one month of data taking)

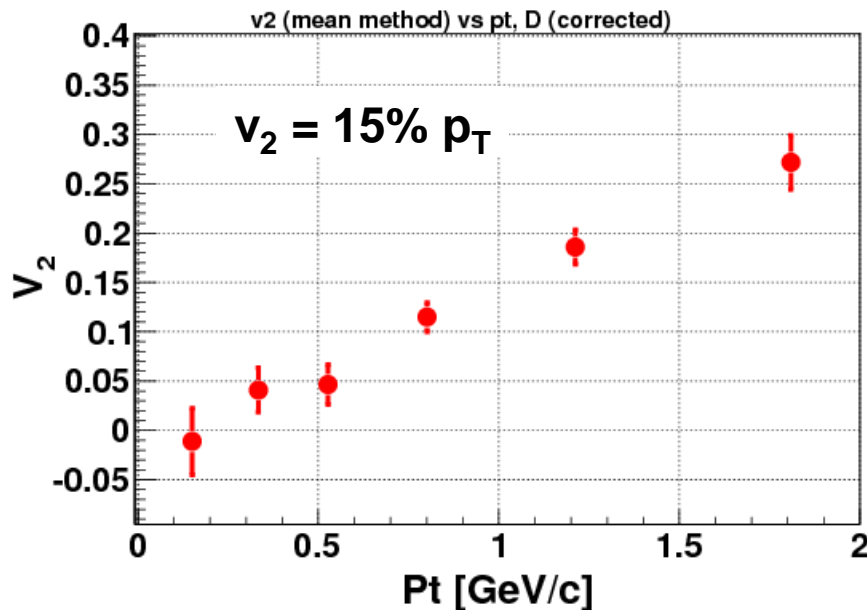


- **A** quite good accuracy on the integrated v_2 is obtained, even for a moderate elliptic flow magnitude
- **T**here is still room for improving the reaction plane resolution, which would reflect in a significantly more accurate measurement

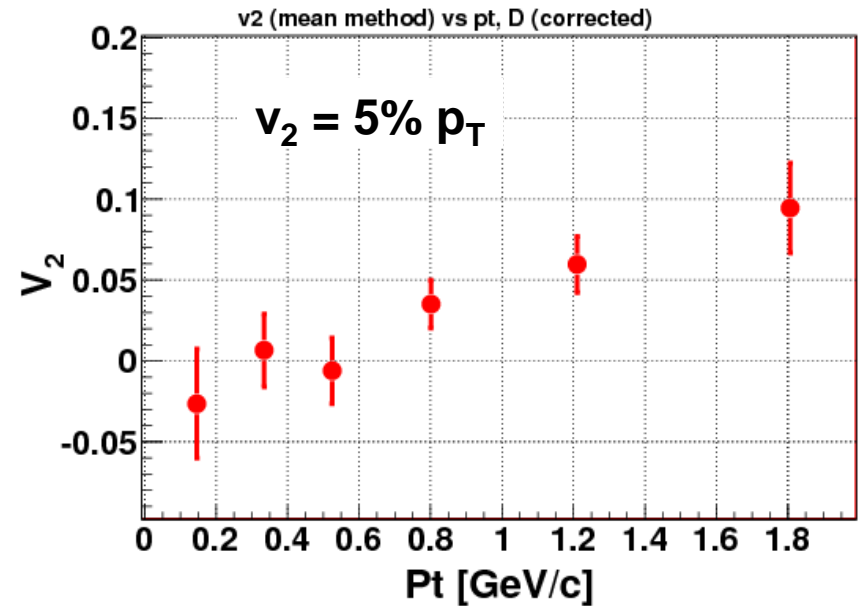
v_2 reconstruction for D-Mesons

- Differential $v_2(p_T)$ -

- We assume 50K reconstructed D-mesons (D^+ , D^- , D^0 , D^0 bar, etc) (\sim one month of data taking)
- A reaction plane resolution of 40 degrees



Strong charm elliptic flow
→ small statistical errors



Moderate charm elliptic flow
→ quite large statistical errors

Still, the two scenarios can be distinguished

Summary and Conclusion

Elliptic flow is one of the most promising probes of deconfinement transition, but is also one of the most challenging differential analysis at FAIR energy regime.

A simulation of the reaction plane reconstruction has been performed with the CBM set-up and a fairly good accuracy has been found for semi-peripheral collisions.

A first indicative feasibility study of open charm elliptic flow reconstruction has been conducted, and concluded that the integrated v_2 of individual species will be measurable with a good statistical precision, already after 1 month of data taking, even for a moderate v_2 magnitude.

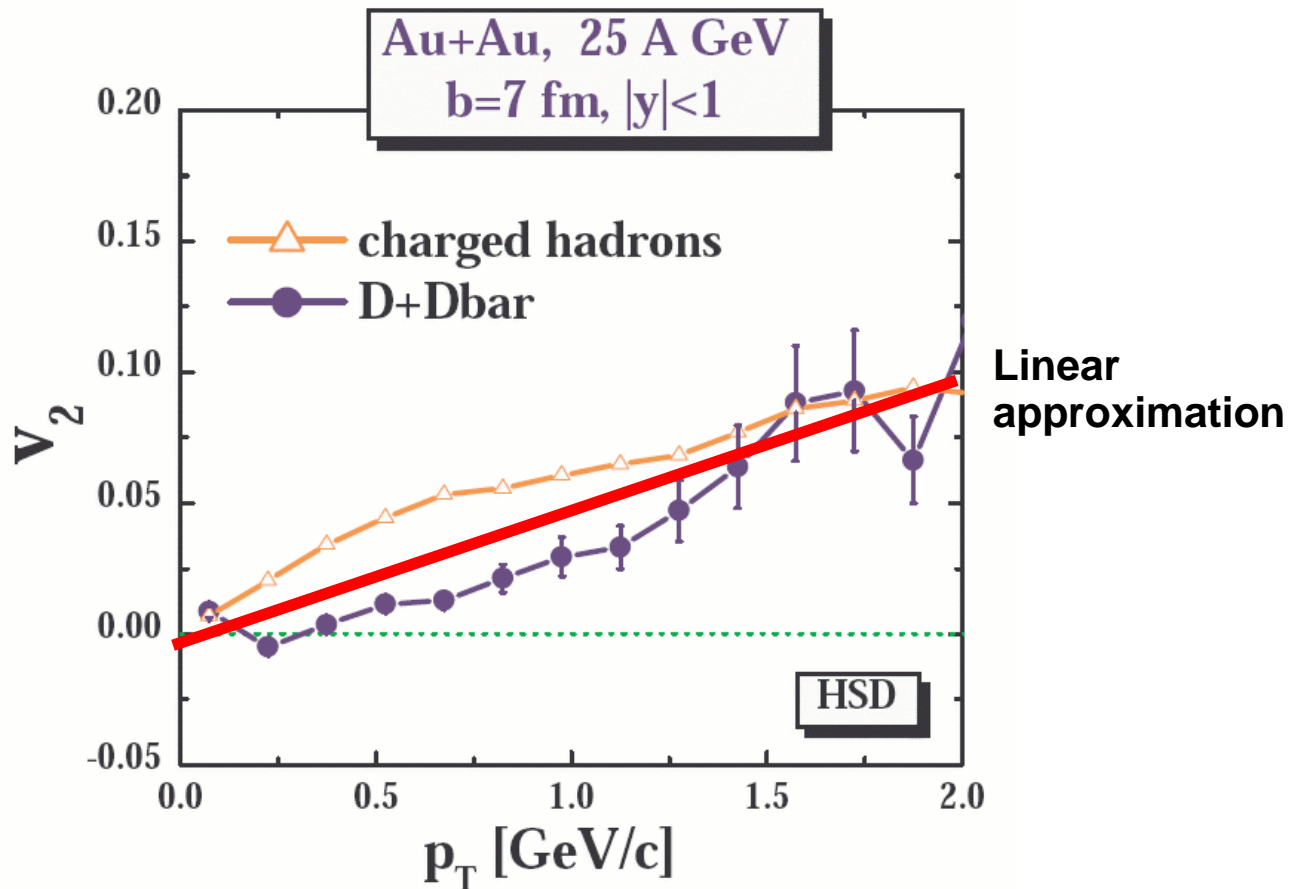
The differential elliptic flow measurement for D-mesons (including all species) has been found to be feasible with a reasonable statistical precision, after only 1 month of data taking.

A longer running period would allow a detailed study of individual species.

Back-up

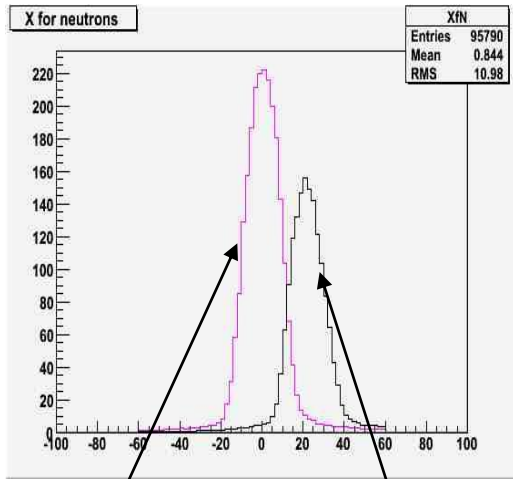
V_2 reconstruction for the D^+

- V_2 (pt) prediction from HSD at FAIR energies -



$V_2^{\text{true}} = 0.05 \times p_T$ for D^+ mesons is assumed for the V_2 reconstruction feasibility study

PSD centered at $X=8.9$ cm (pass 25 GeV beam)



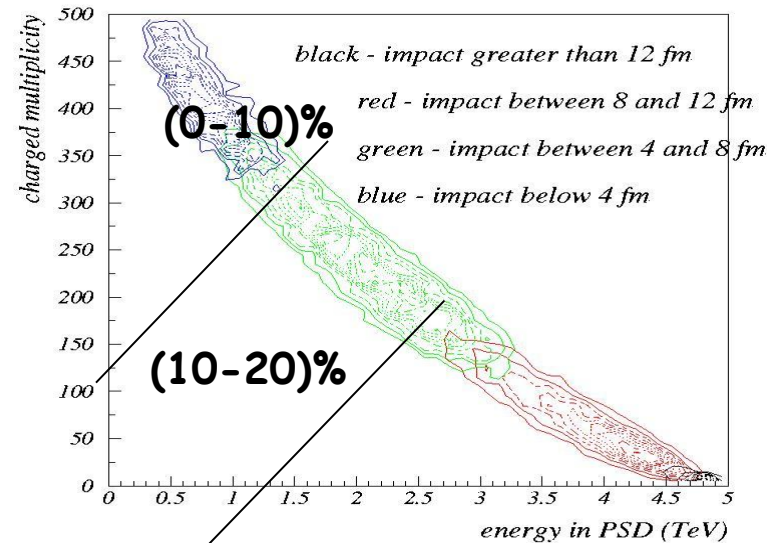
neutrons

protons

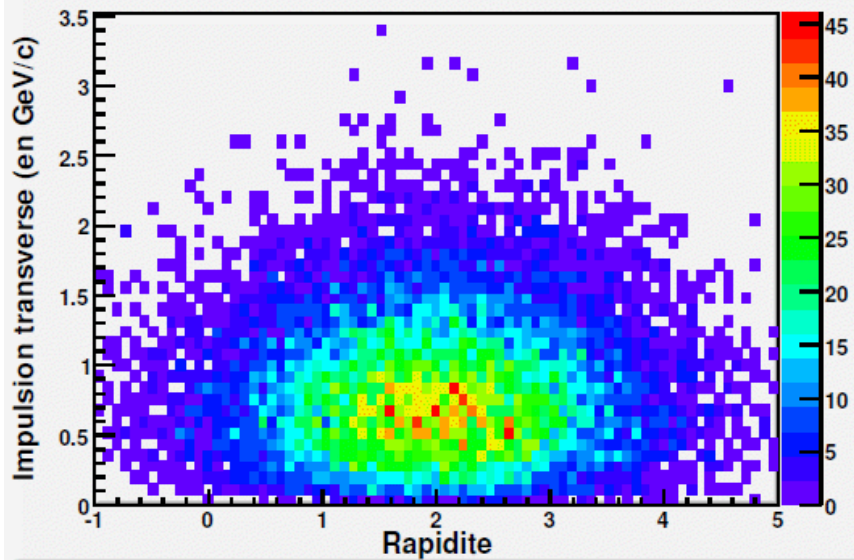
radius

mean energy in circle band (GeV)

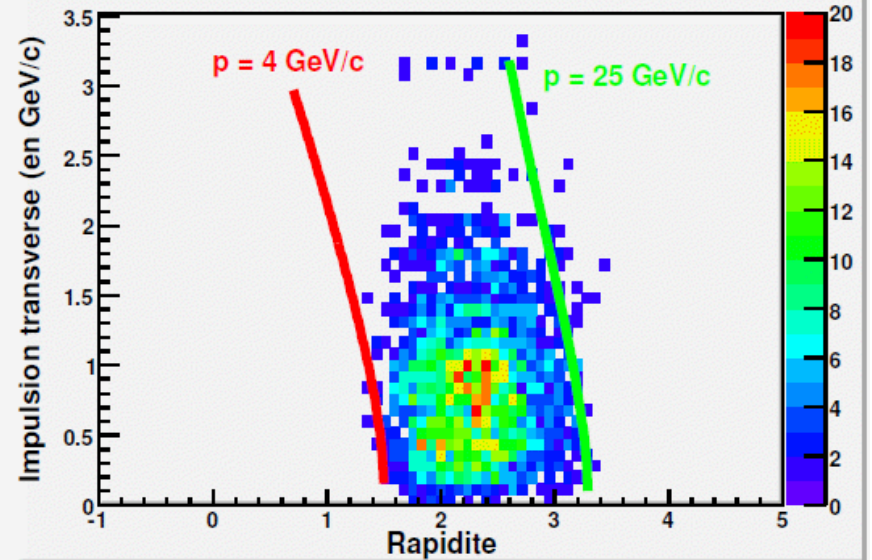
radius	mean energy in circle band (GeV)
2.5 – 7.5	330 (4-8fm) 680 (8-12fm)
7.5 – 12.5	520 (1040)
12.5 – 17.5	540 (1040)
17.5 – 22.5	410 (710)
22.5 – 27.5	260 (340)



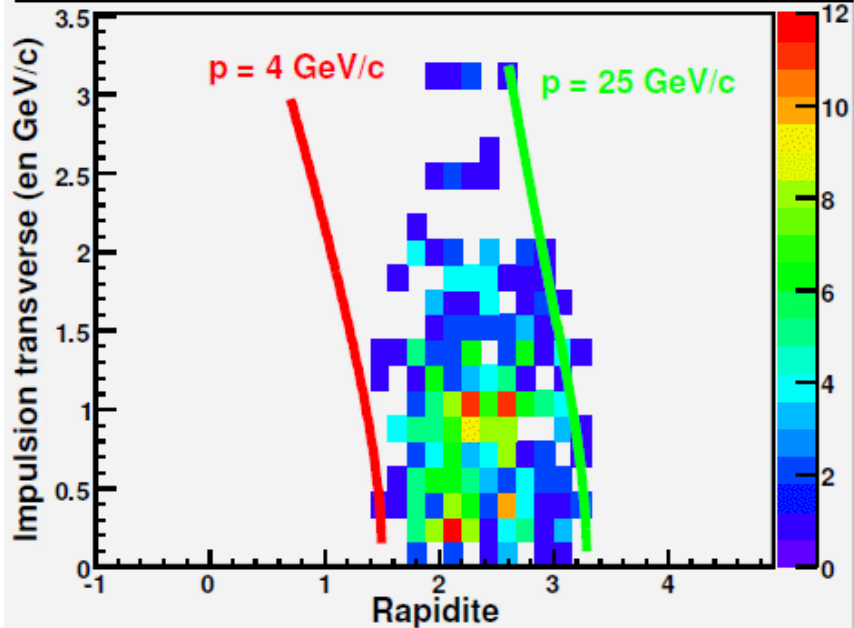
Espace des phases des mesons D+ generes



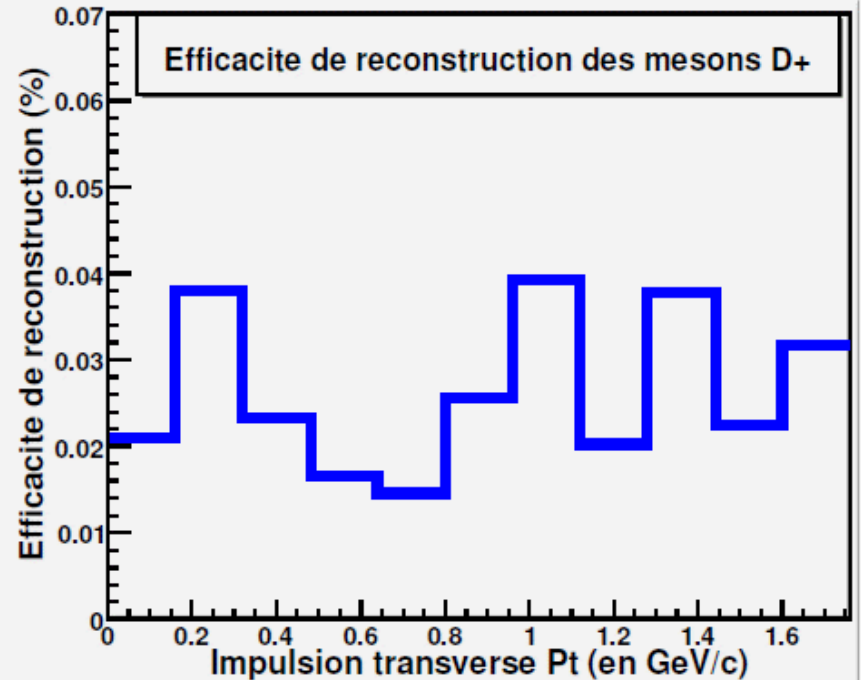
Acceptance geometrique des mesons D+

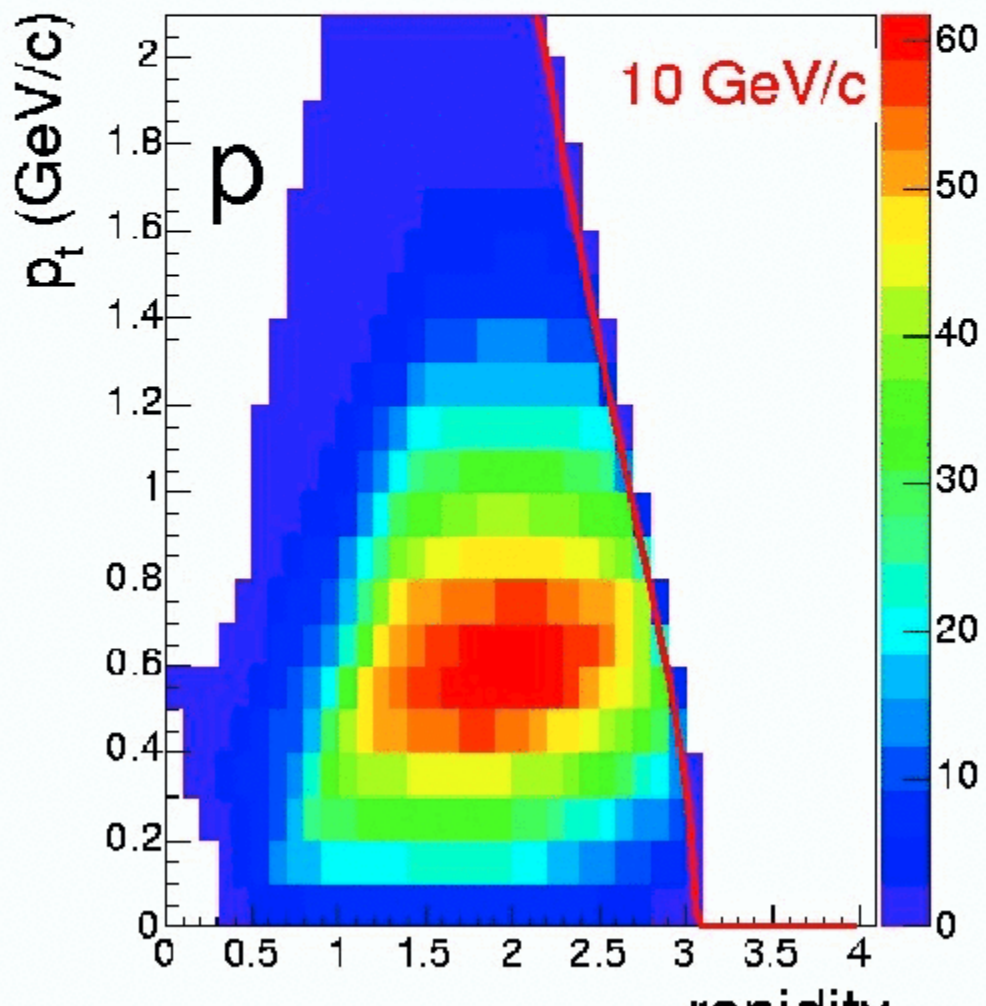


Acceptance des mesons D+ apres la selection

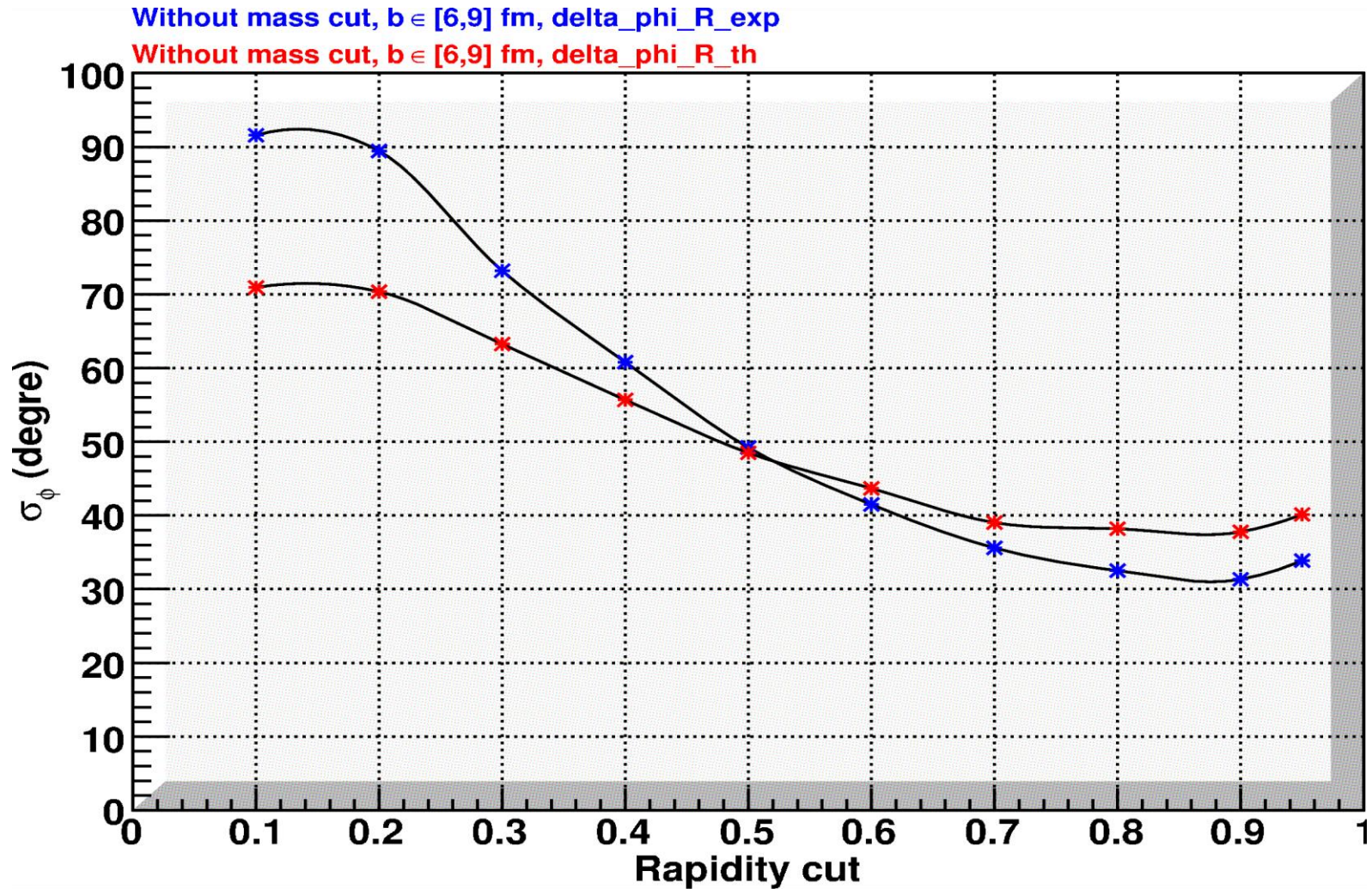


Efficacite de reconstruction des mesons D+





Results : σ_ϕ versus rapidity cut - theoretical and experimental



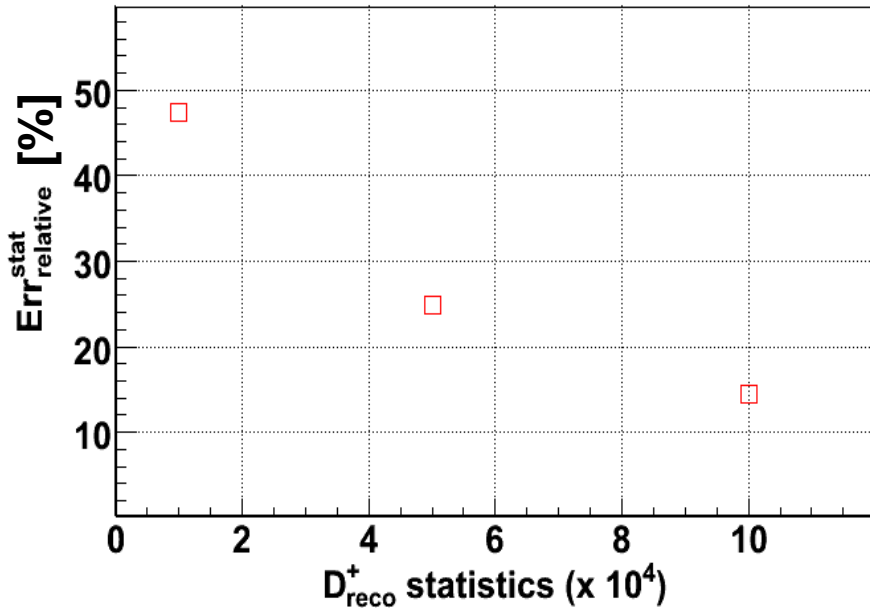
σ_{exp} : Ψ_{RP}^{reco} from 2 sub-events with equal multiplicity : $(\Psi_{RP}^{reco1} - \Psi_{RP}^{reco2}) / 2 \sim \sigma_\phi$
 if σ_ϕ relatively small !!

V_2 reconstruction for the D^+

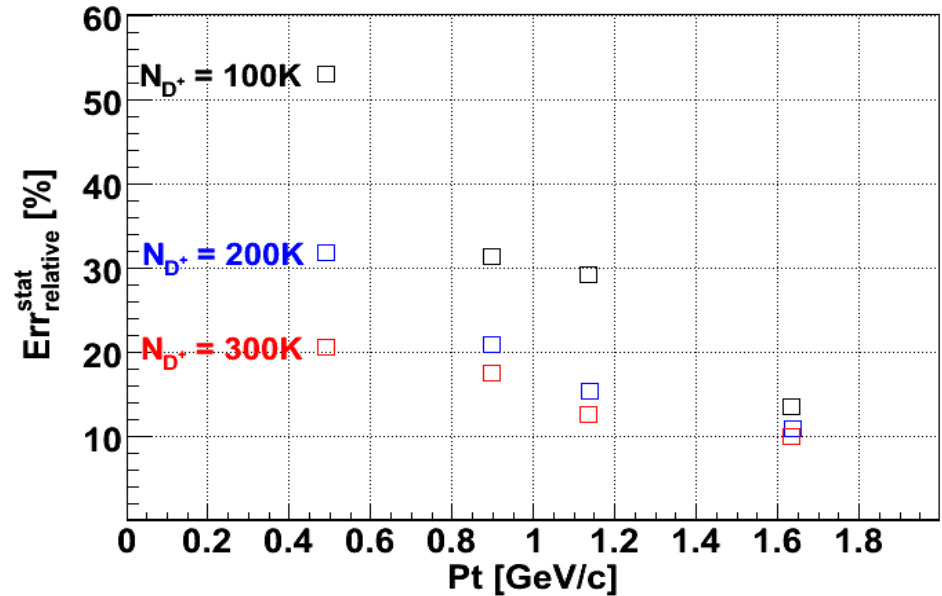
- Results -

Integrated and differential V_2 (Pt) for $V_2 = 0.05 \times \text{Pt}$ and $\sigma_{\text{RP}} = 40$ degrees

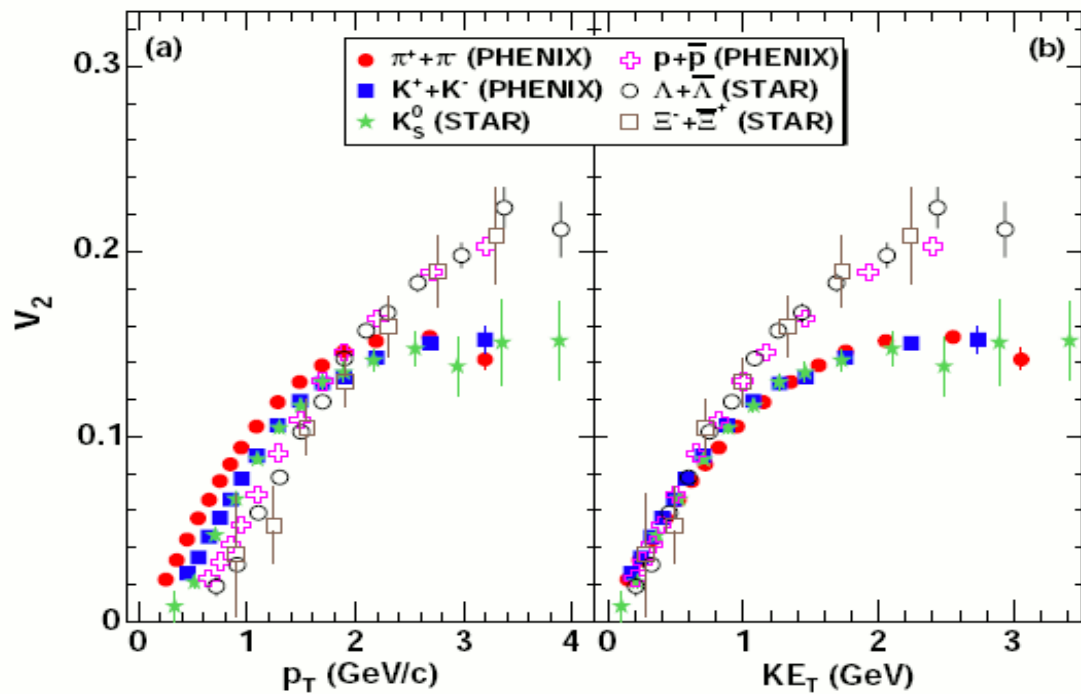
Fit method, full pt range



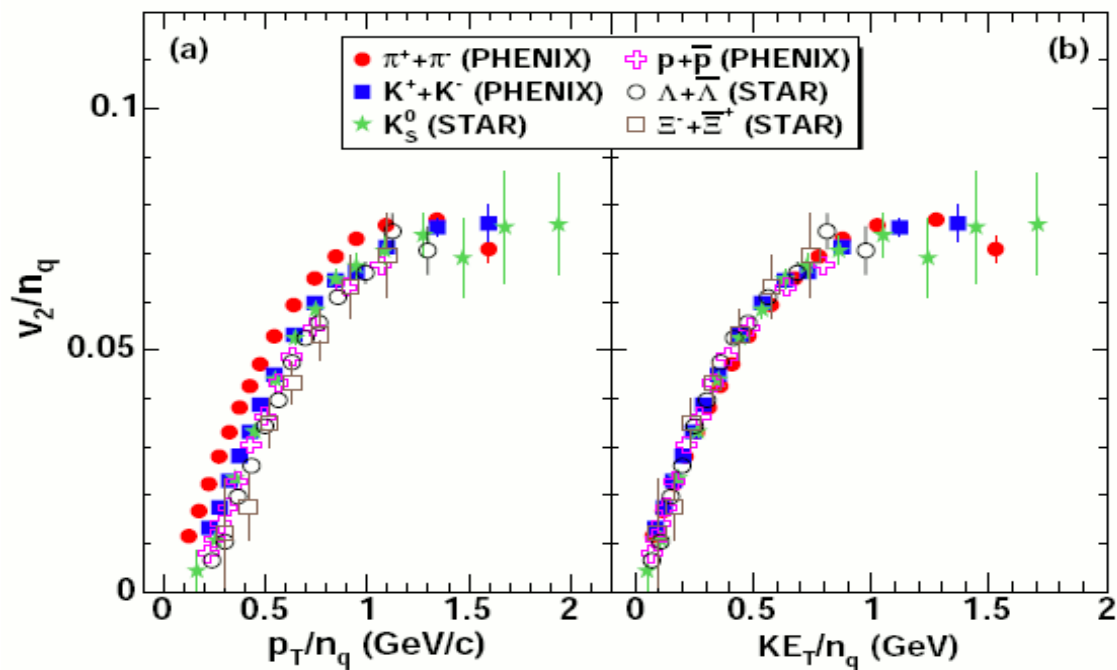
Need ~ 70 000 D^+



Need ~ 300 000 D^+

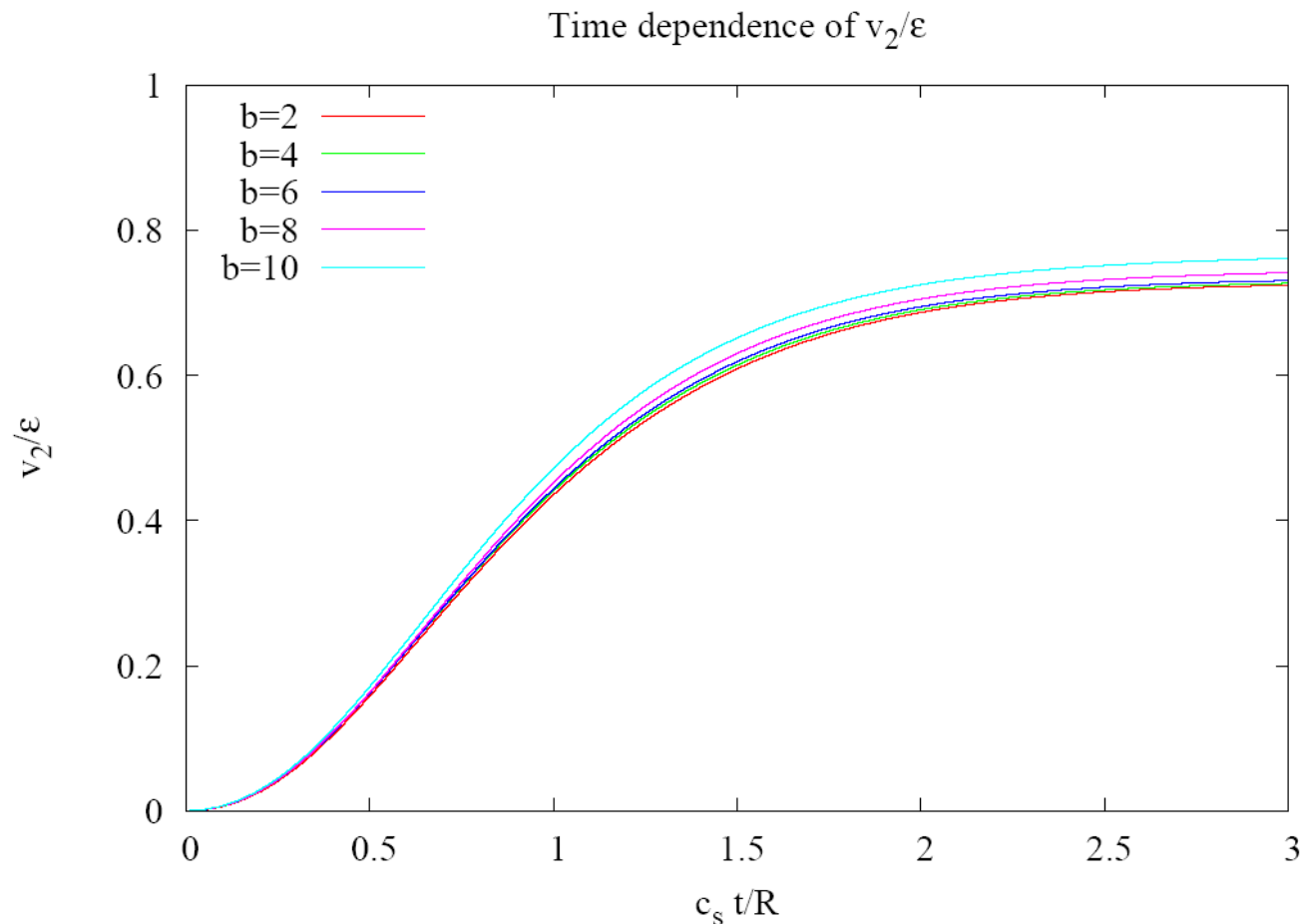


arXiv:nucl-ex/0608033v1

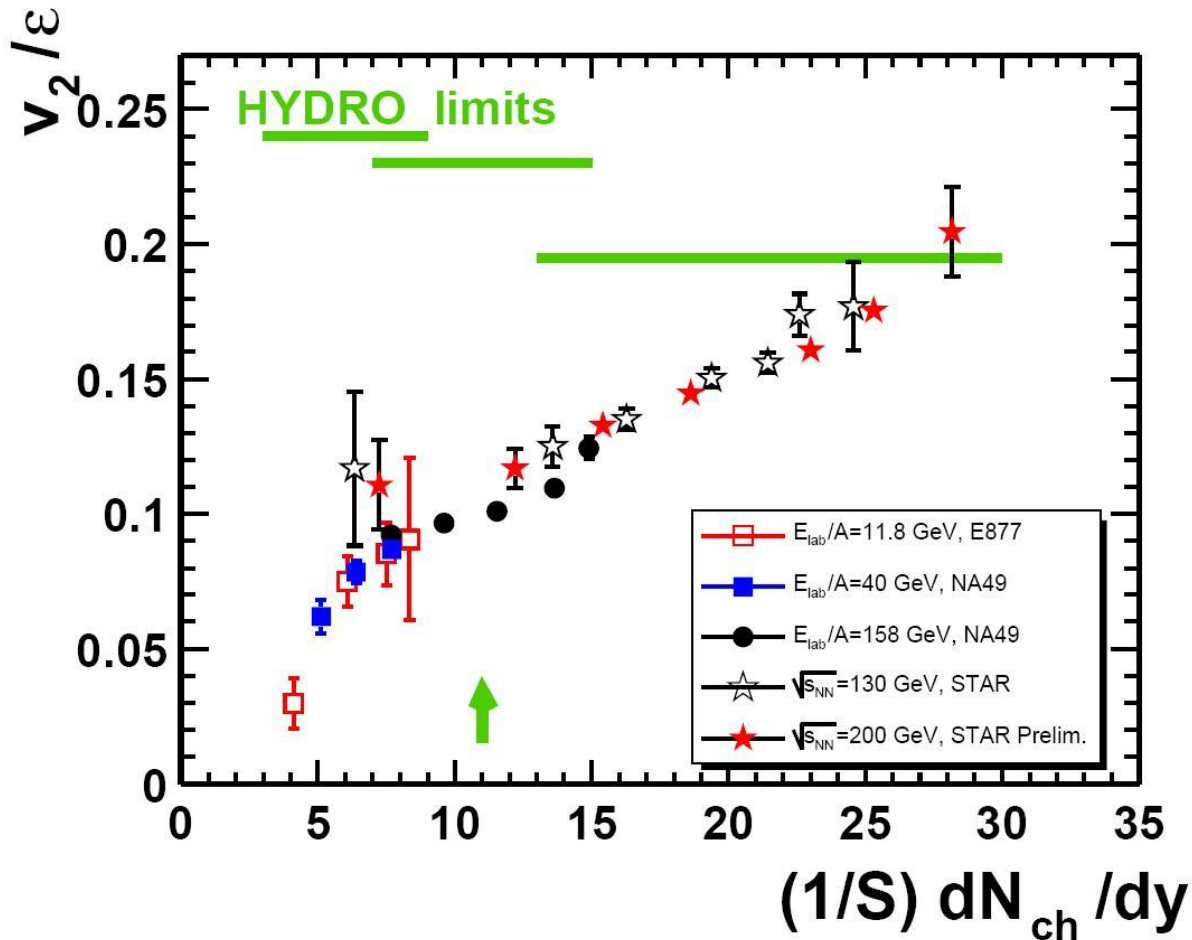


1. When is v_2 created?

At a time $\sim R/c_s$, where $R=(1/\langle x^2 \rangle + 1/\langle y^2 \rangle)^{-1/2}$



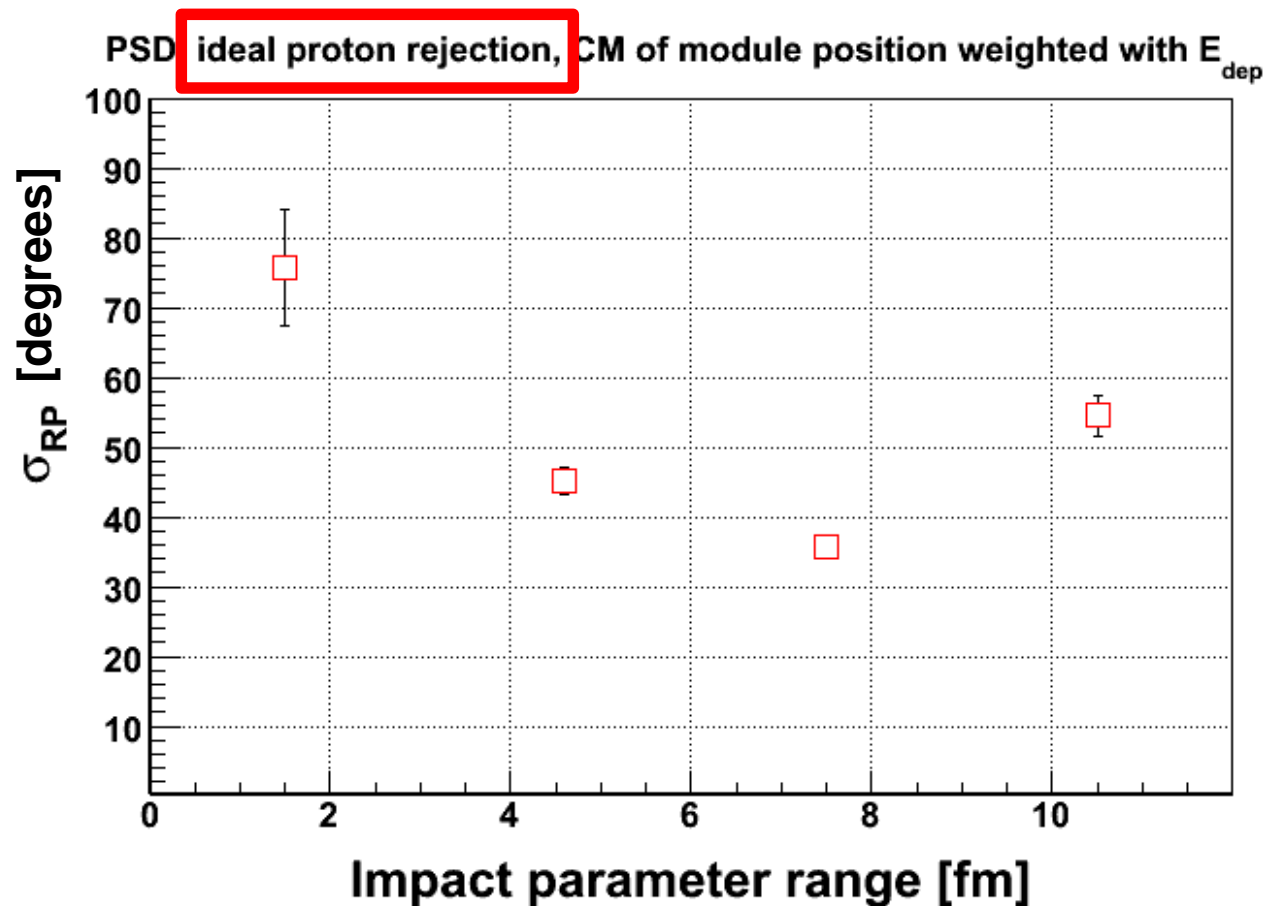
v_2/ϵ : Data from SPS and RHIC



Hydrodynamical evolution at RHIC energy regime square(s) ~ 200 GeV

Reconstruction of the reaction plane

- Systematic study -



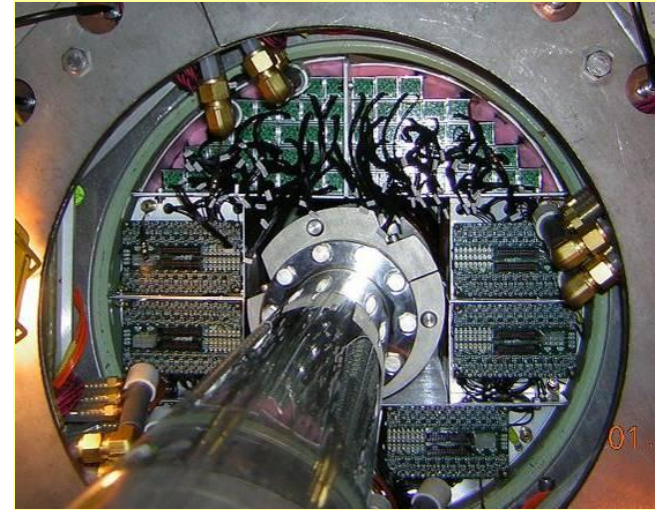
The best reaction plane resolution is obtained for b between 6 and 9 fm
 $\sigma_{\text{RP}} \sim 40$ degrees is used for the feasibility study of v_2 reconstruction

Reaction Plane Measurement with PHENIX



Beam-Beam Counters (BBC)

- Quartz Cherenkov radiators
- 64 elements in 3 rings
- $3.0 < |\eta| < 4.0$
- All Runs

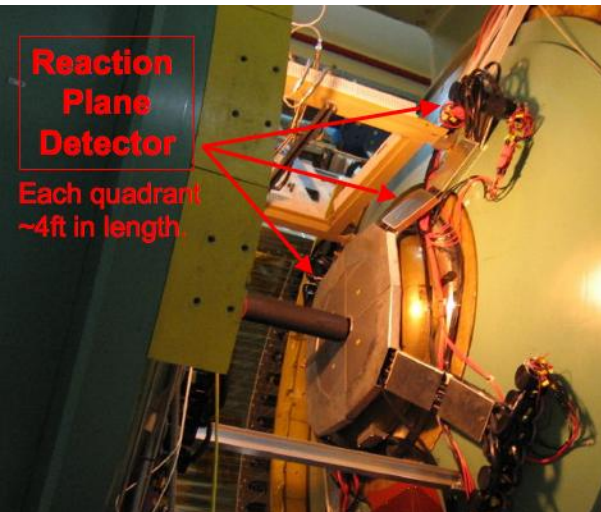


Reaction Plane Detector (RxnP)

- plastic scintillators @ $38 < |z| < 40$ cm
- 12 segments in ϕ
- 2 segments in η
 - $1.0 < |\eta| < 1.5$
 - $1.5 < |\eta| < 2.8$
- Pb converter
- Run 7+

Muon Piston Calorimeter (MPC)

- PbWO_4 PHOS crystals
- 192 towers
- $3.1 < |\eta| < 3.7$
- Run 6+



Multiple overlapping and complementary measurements

Starting from 2007, PHENIX uses two new Reaction Plane Detectors (RxnP) (see Fig. 1) to measure the reaction plane of each collision following methods described in [13, 14]. It improves the reaction plane resolution, and thus gives a correction $\sigma_{RP} = \langle \cos(2\Delta\psi_{RP}) \rangle$ twice better than what was achieved previously with the Beam Beam Counters (BBC), or what can be measured with the Muon Piston Calorimeters (MPC) (see Fig. 2).

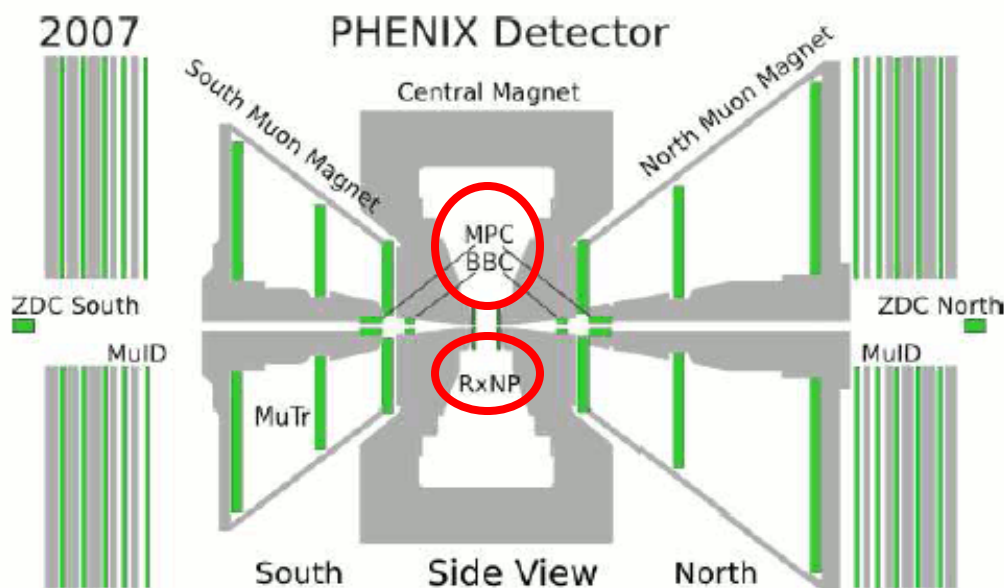


Figure 1. PHENIX detector during 2007 data taking with the RxnP detector near the collision vertex.

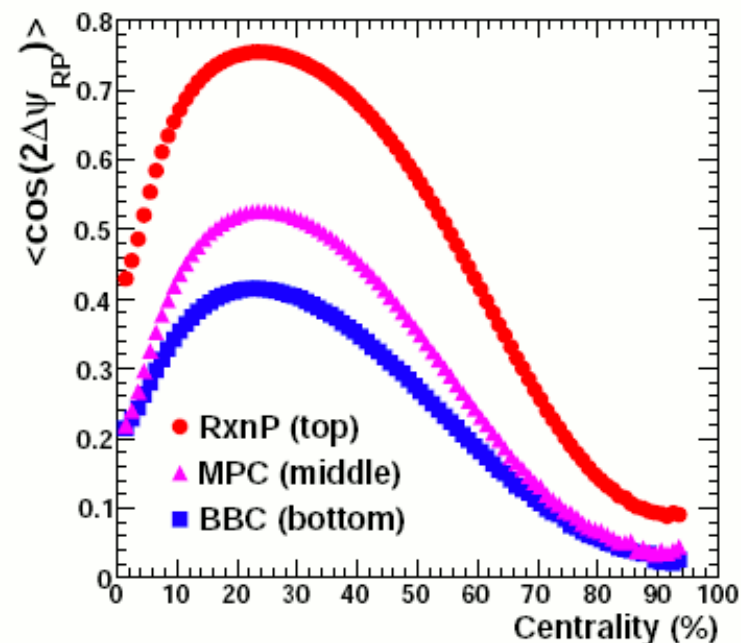


Figure 2. Reaction plane resolution as a function of centrality, measured with the RxnP (top), the MPC (middle), or the BBC (squares).

“Shifting Method” in PHENIX

$$\Psi'_1 = \Psi_1 + \Delta\Psi_1.$$

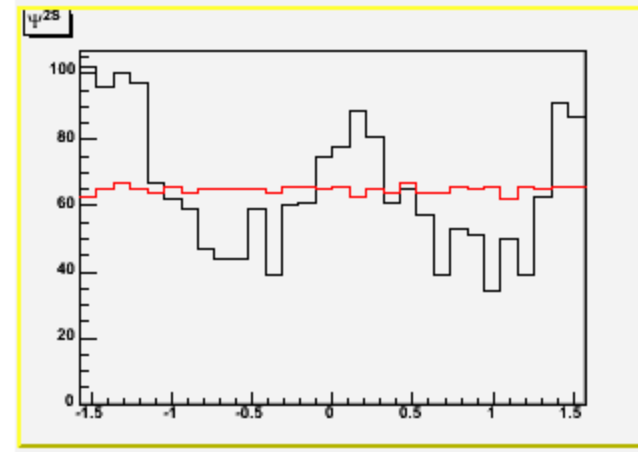
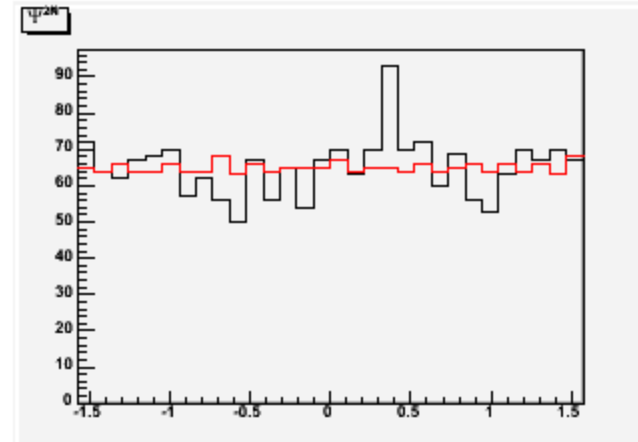
$$\Delta\Psi_1 = \sum_n (A_n \cos(n\Psi_1) + B_n \sin(n\Psi_1))$$

$$B_n = \frac{2}{n} \langle \cos(n\Psi_1) \rangle,$$

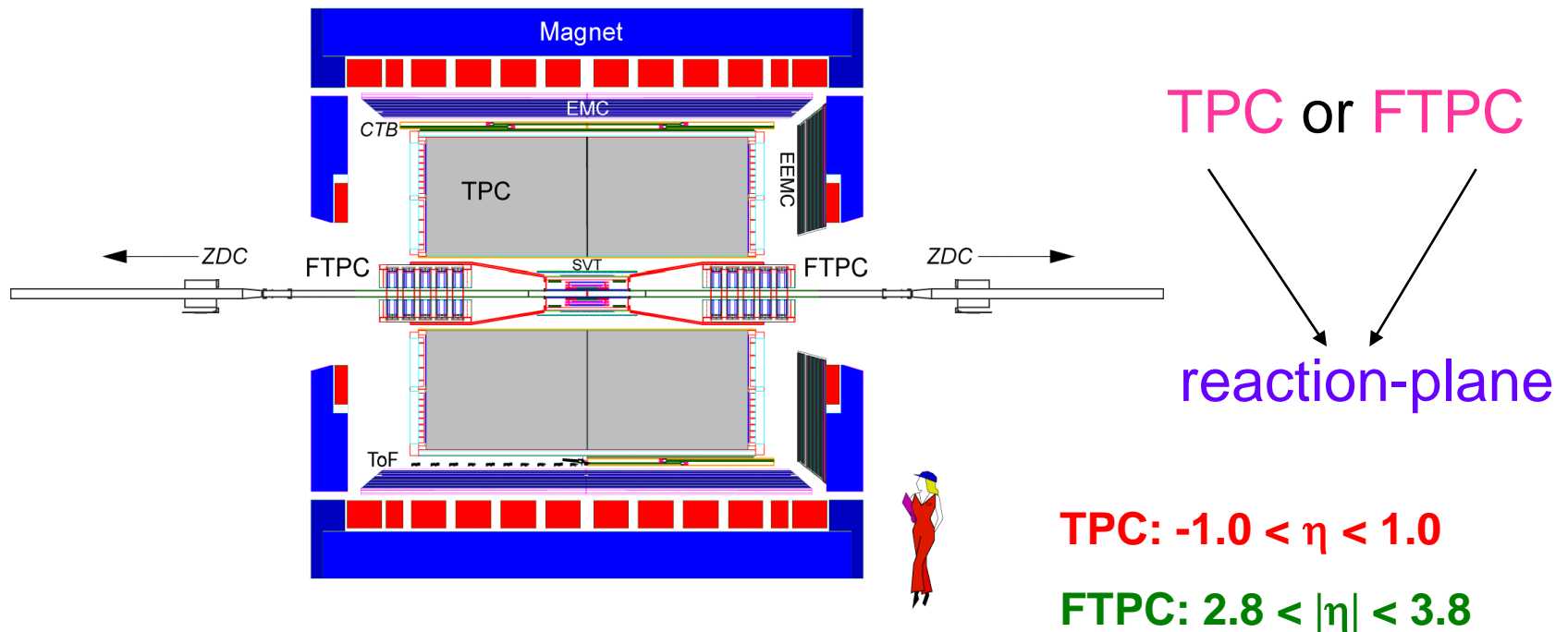
$$A_n = -\frac{2}{n} \langle \sin(n\Psi_1) \rangle,$$

Usually $1 \leq n \leq 32$

$$\Psi'_1 = \Psi_1 + \sum_n \frac{2}{n} (-\langle \sin(n\Psi_1) \rangle \cos(n\Psi_1) + \langle \cos(n\Psi_1) \rangle \sin(n\Psi_1))$$



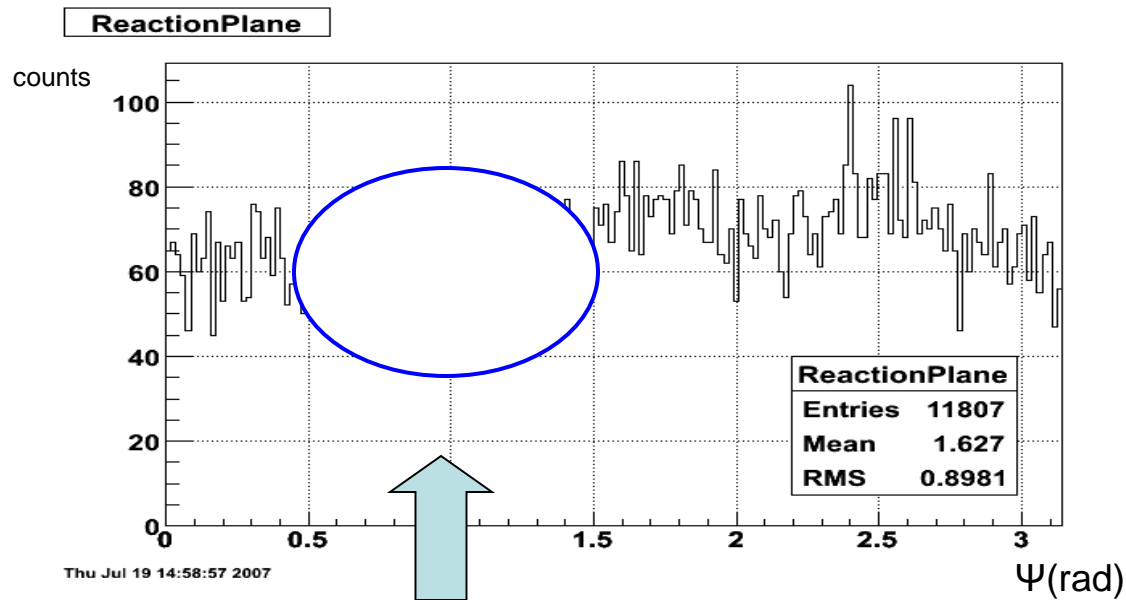
Reaction Plane Determination in STAR



Why we use FTFC reaction-plane?

1. The non-flow correlation and auto correlation effects are smaller from FTFC reaction-plane
2. Comparing the v_2 results from FTFC reaction-plane with TPC reaction-plane can help us understand the non-flow effects.

Shifting Corrections in STAR

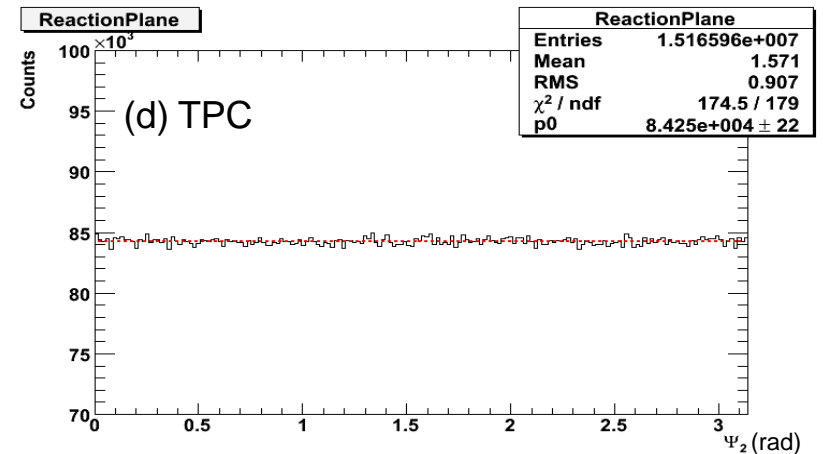
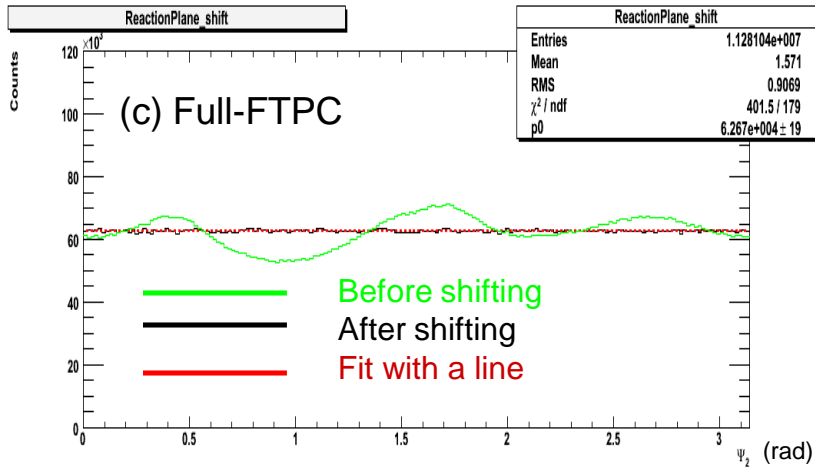
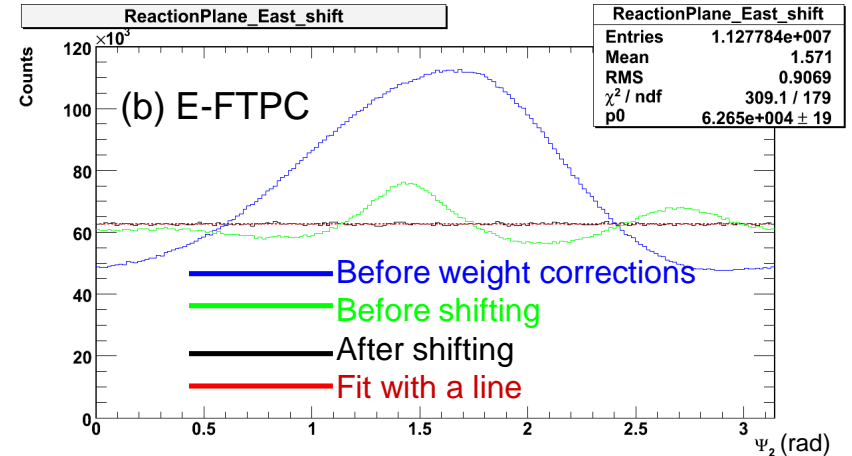
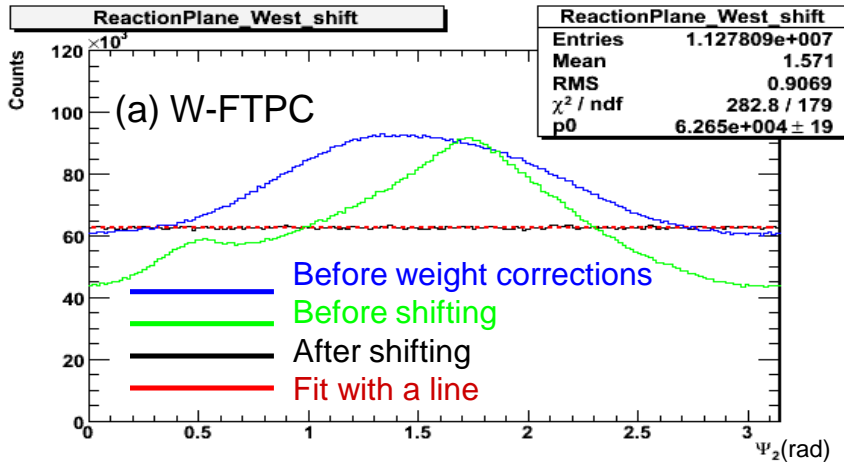


Even with the weighting method, the event plane distribution from FTPC we obtain isn't isotropic. We can flatten that distribution by shifting:

$$\Psi' = \Psi + \sum_n \frac{1}{n} (-\langle \sin 2n\Psi \rangle \cos 2n\Psi + \langle \cos 2n\Psi \rangle \sin 2n\Psi)$$

J. Barrette et al. Phys.Rev. C56, 3254(1997), nucl-ex/9707002

Reaction Plane from FTPC and TPC inSTAR

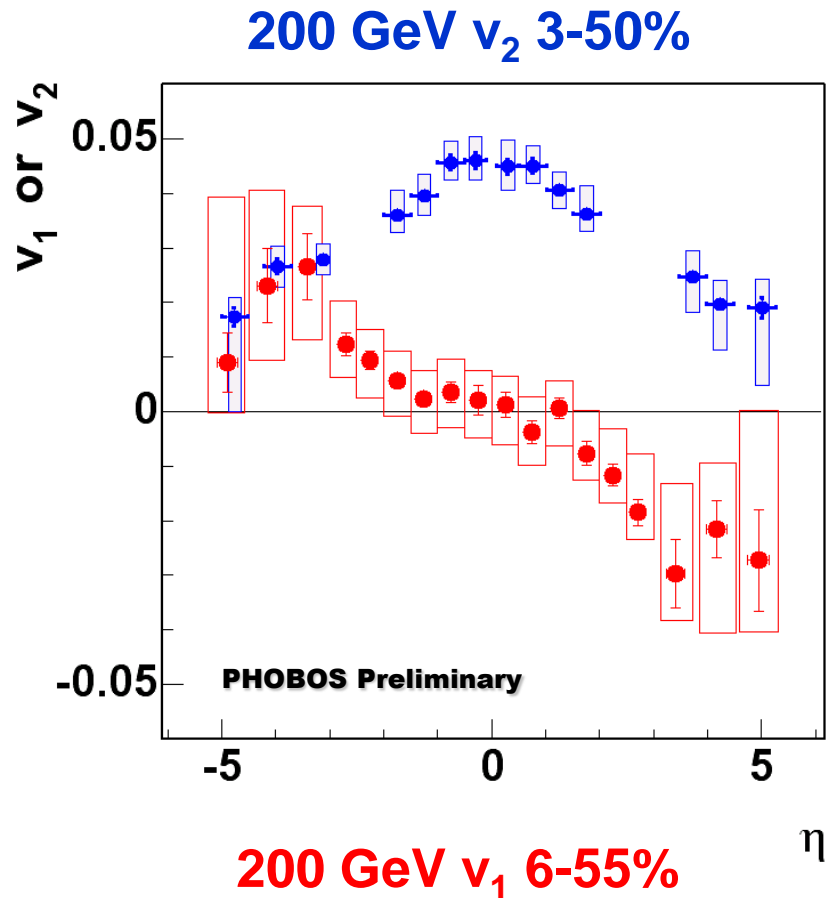


Full-FTPC: $\cos(2\Delta\Psi): 0.335$

TPC: $\cos(2\Delta\Psi): 0.604$

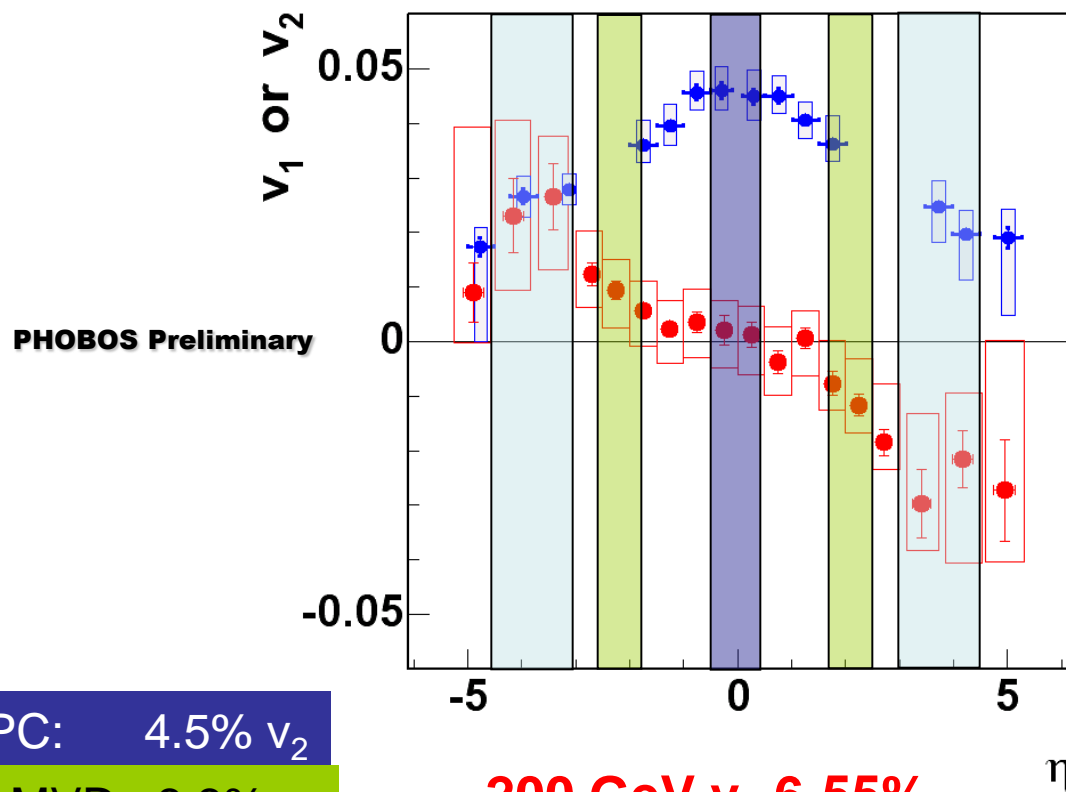
How much asymmetry?

- v_1 , v_2 measure asymmetry of system
- Forget flow for now
- effects are at the few % level
 - nontrivial measurement
 - must understand asymmetric efficiency/ acceptance/ background



How much asymmetry? ctd

200 GeV v_2 3-50%



v_2 decreases with η
 v_1 increases with η

MVD sees somewhat stronger v_2 signal than BBC

STAR TPC: 4.5% v_2

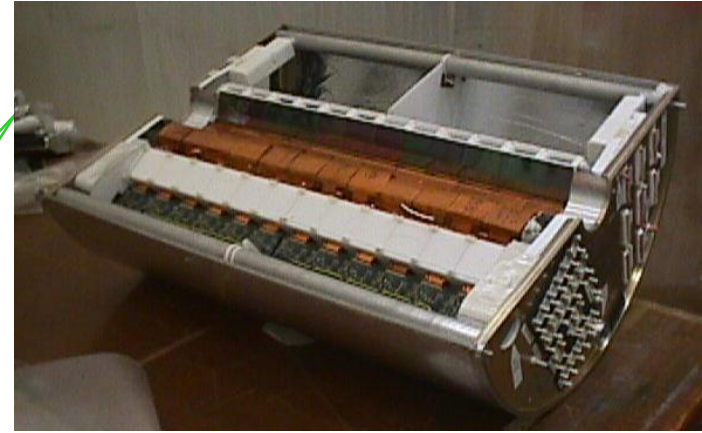
PHENIX MVD: 3.3% v_2

PHENIX BBC: 2.8% v_2

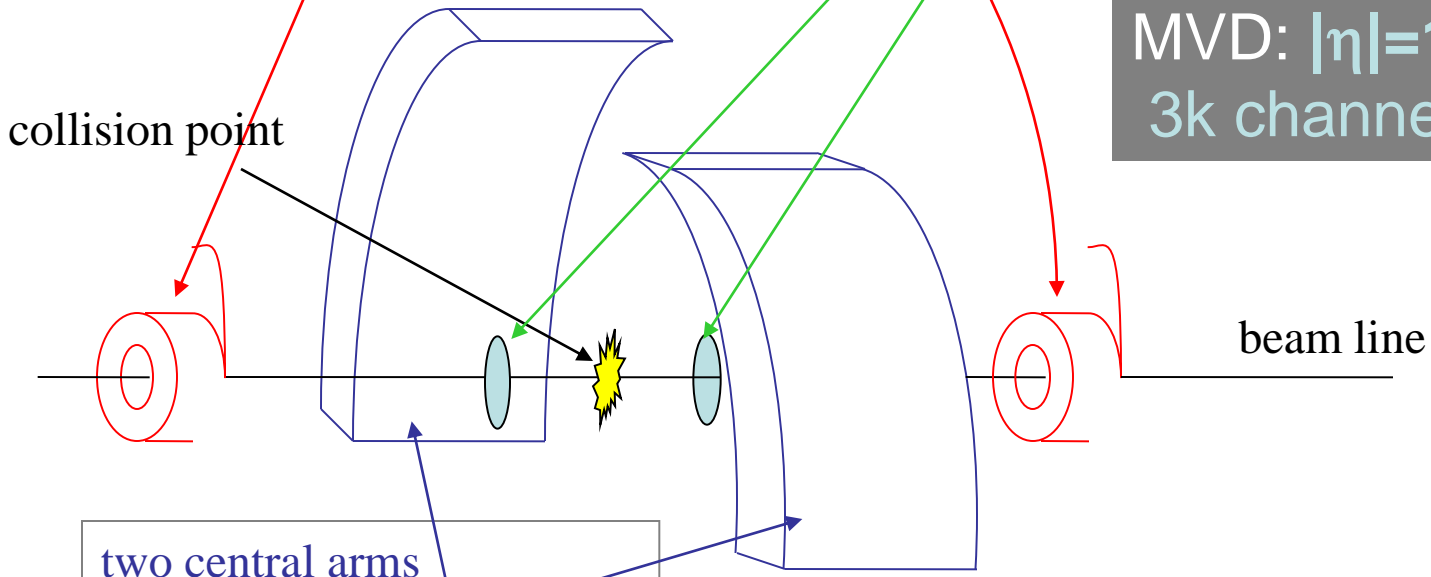
BBC and MVD in PHENIX



Beam-beam counter (BBC) $|\eta|=3\sim 4$
64pmts in each BBC
charged particles



MVD: $|\eta|=1.8\sim 2.6$
3k channels per side!



collision point

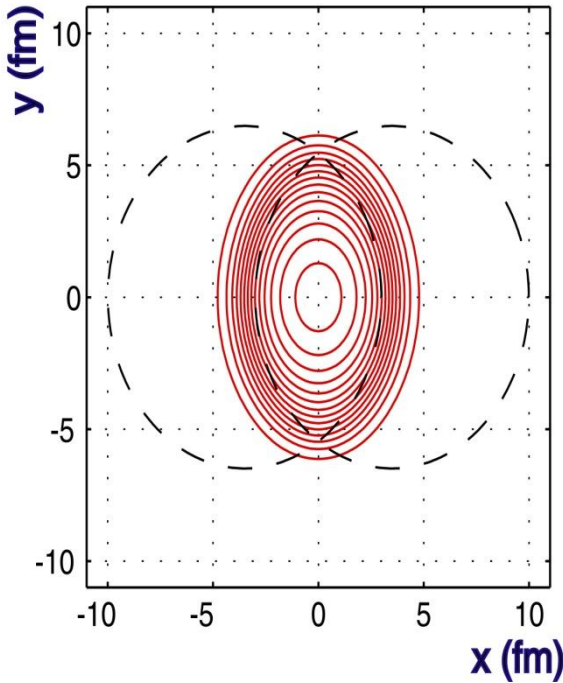
beam line

two central arms (CNT) $|\eta|<0.35$
Dch,PCs,TOF,EMCAL
tracking, momentum, PID

ion and Forward Upgrades
Workshop
Santa Fe, NM

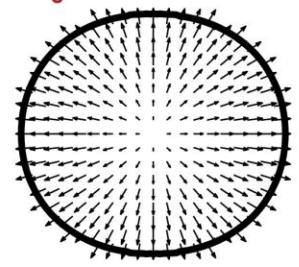
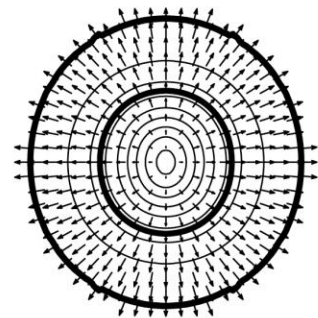
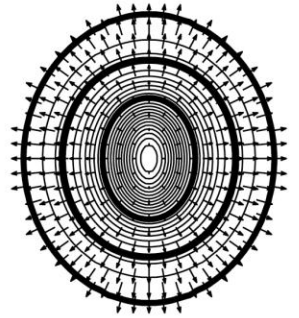
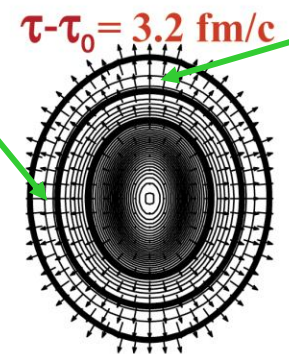
Time evolution in a ideal hydrodynamic model calculation

$Pb + Pb, b = 7 \text{ fm}$



dn/dr (dP/dr) higher

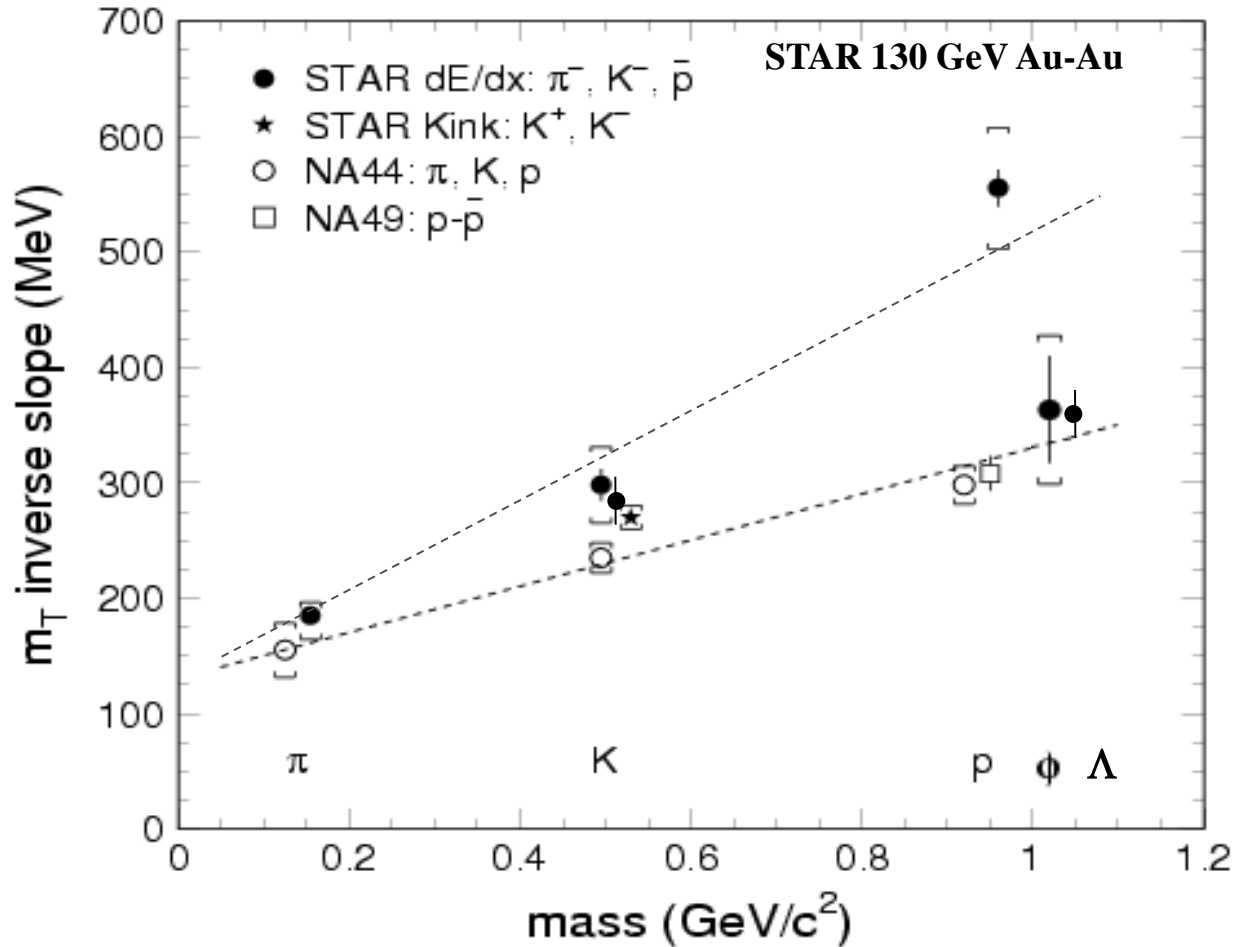
dn/dr (dP/dr) lower



- Elliptic Flow reduces spatial anisotropy -> shuts itself off

$PV = ART \rightarrow dP \sim dn$

Mass dependence of m_T slope - Radial flow



- The different m_T slopes indicate that the slope should not be interpreted as a simple thermal temperature.

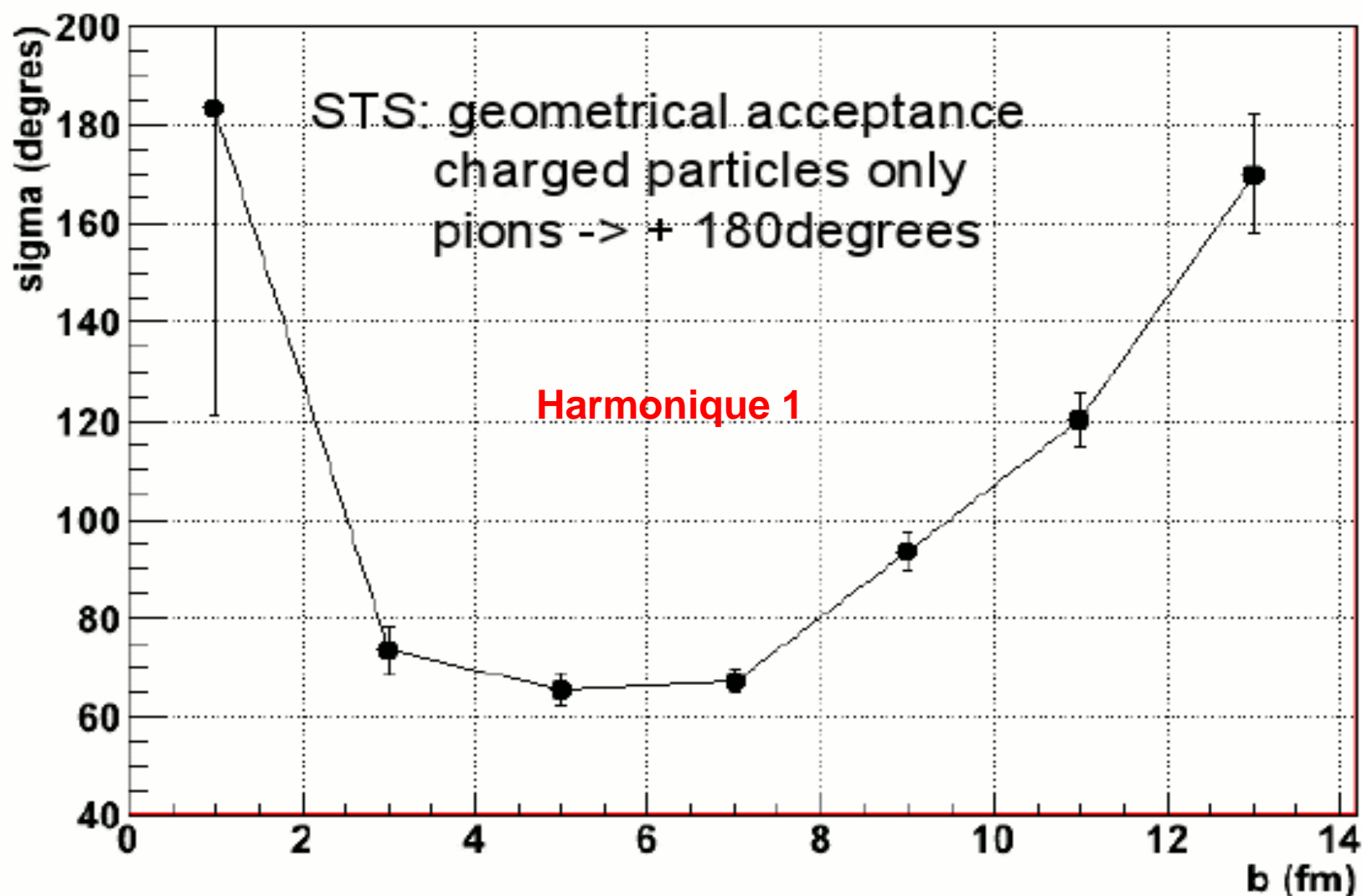
- It also has a component from **collective radial velocity**.

$$\frac{1}{m_T} \frac{d^2 N}{dm_T dy} e^{-\frac{m_T - m}{T}}$$

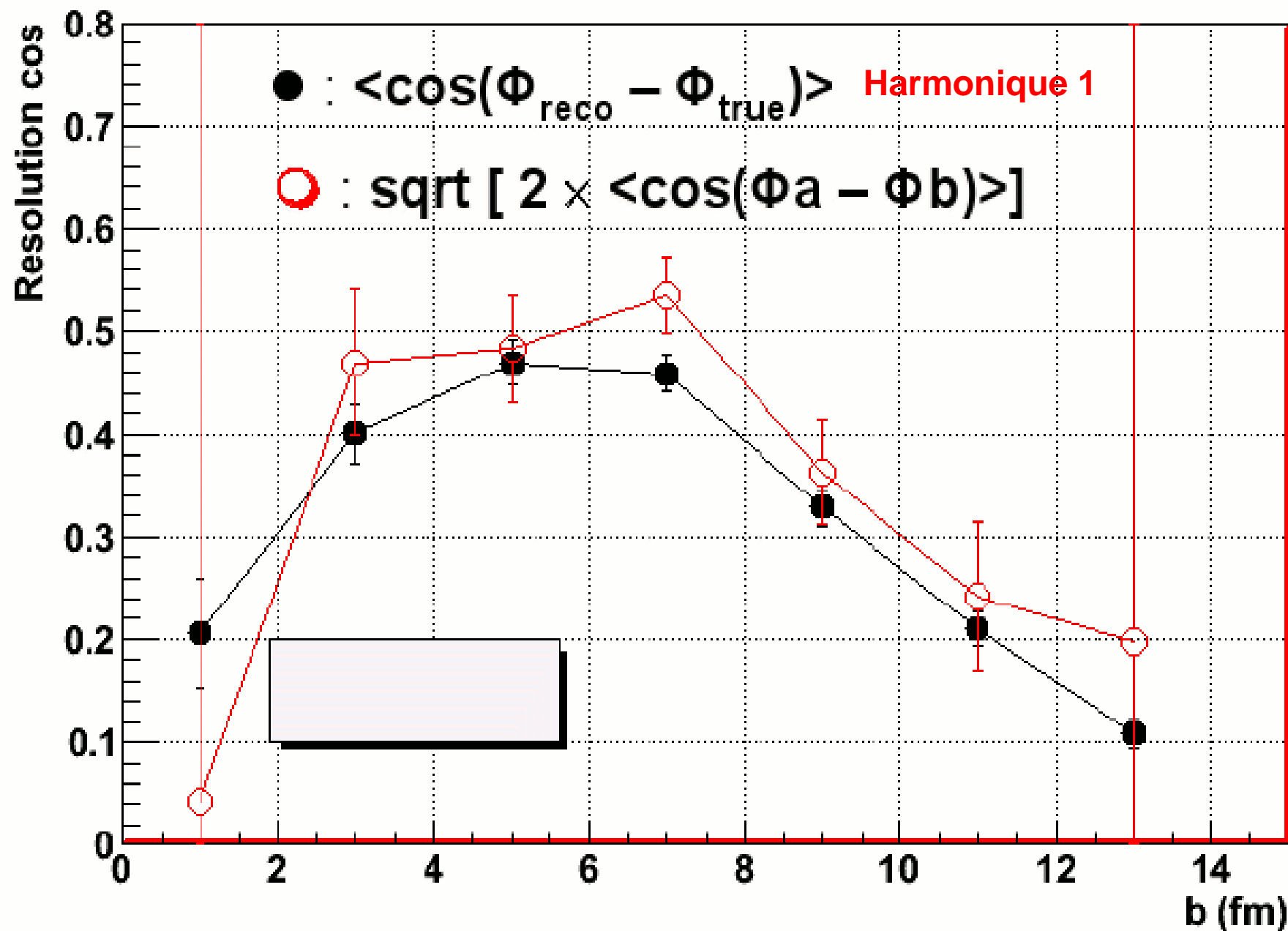
$$T = T_{Th} + m\beta^2$$

Resolution expressed in terms of $\sigma(\Phi_{\text{reco}} - \Phi_{\text{true}})$

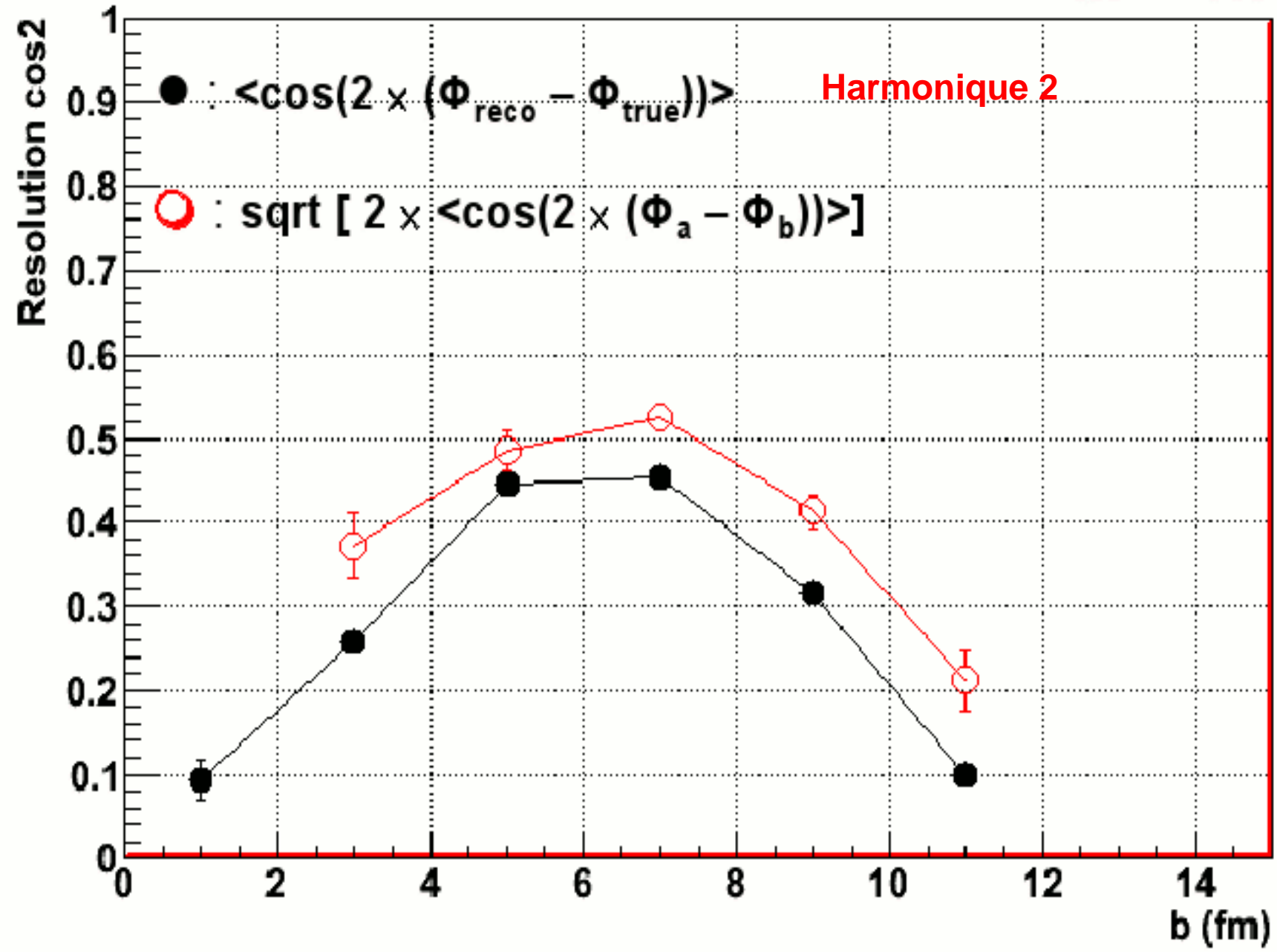
ecart type en fonction du parametre d impact



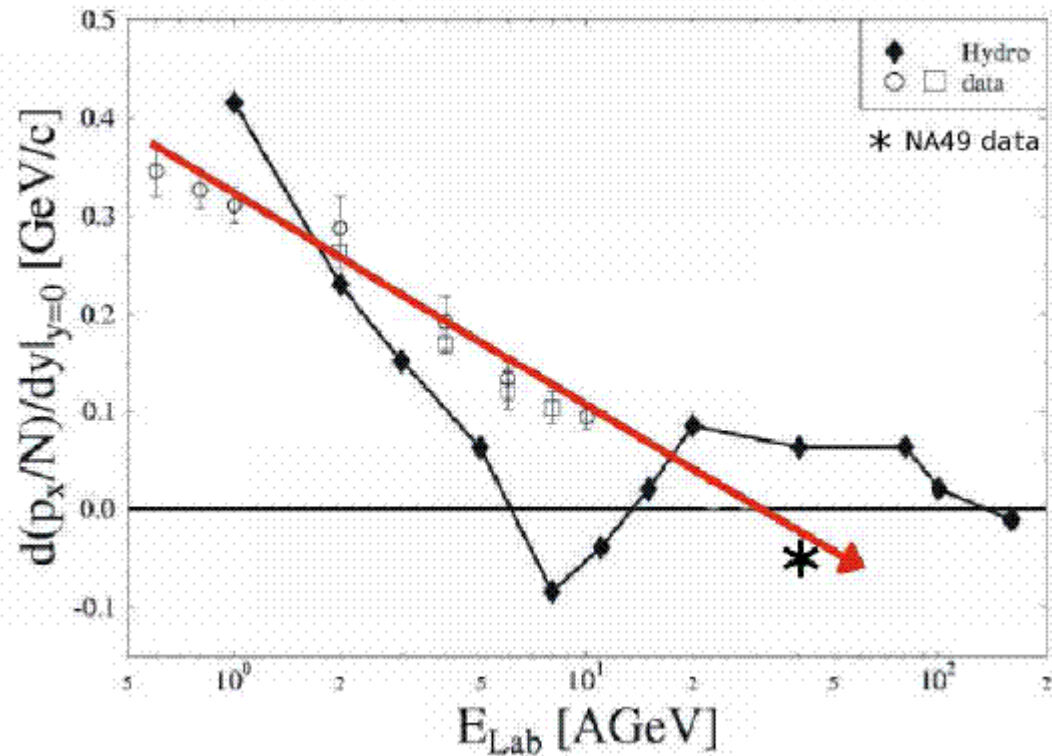
Resolution expressed in terms of $\langle \cos(\Phi_{\text{reco}} - \Phi_{\text{true}}) \rangle$



Resolution expressed in terms of $\langle \cos(2 \times (\Phi_{\text{reco}} - \Phi_{\text{true}})) \rangle$



Collapse of proton v_2 as a probe of the 1st order phase transition at FAIR energies (1)

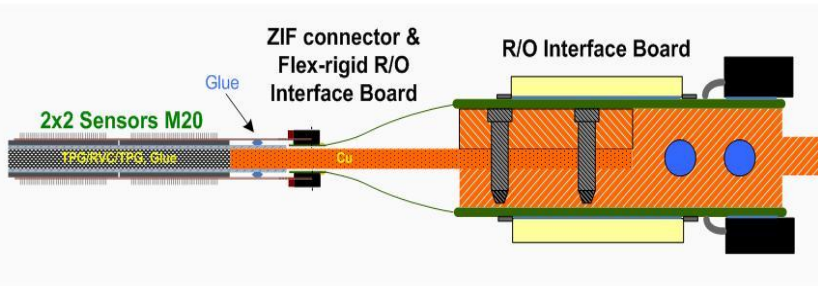


Prediction at FAIR energy (around 30 – 40 A.GeV)

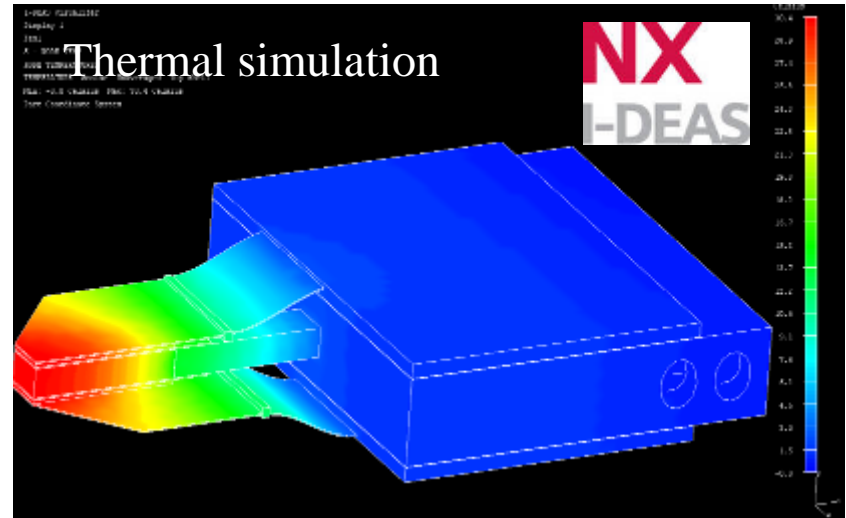
- Collapse of proton v_2
- also: Wiggle of v_1 at mid-rapidity

MAPS sensors for the MVD – integration (2)

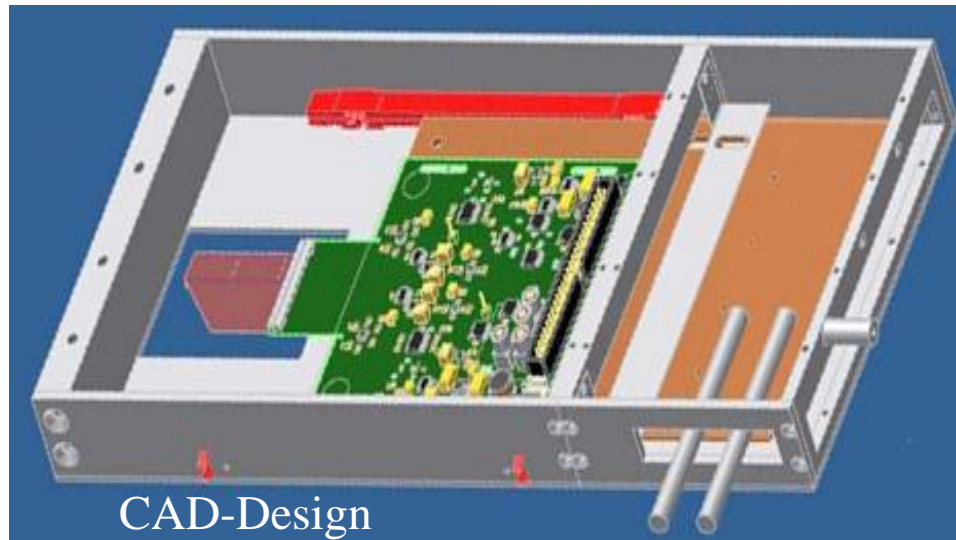
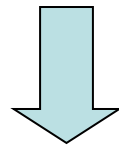
Ladder



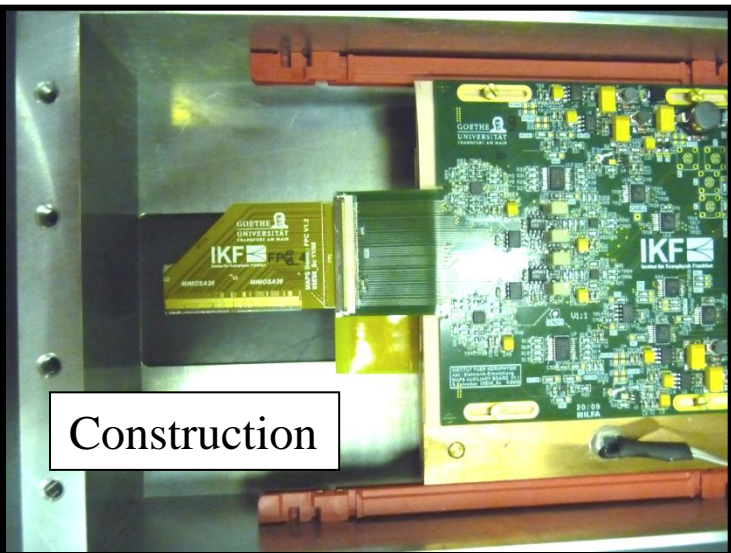
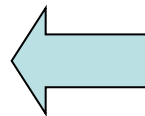
Thermal simulation



Power: $\sim 1\text{ W/cm}^2$



CAD-Design

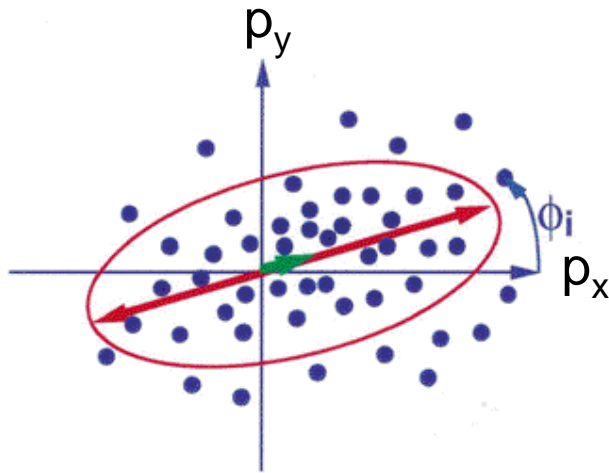


Construction

Flow analysis

Several flow analysis methods: cumulant, Lee-Yang zero method, here: **Event Plane Method**

Transverse Plane



$$\Psi_{\text{plane}} = \tan^{-1} \frac{\sum \sin(\phi_i)}{\sum \cos(\phi_i)}$$

$$2 \Psi_{\text{ellipse}} = \tan^{-1} \frac{\sum \sin(2\phi_i)}{\sum \cos(2\phi_i)}$$

$$Q_n \cos(n\Psi_n) = \sum [w_i \cos(n\phi_i)]$$

$$Q_n \sin(n\Psi_n) = \sum [w_i \sin(n\phi_i)]$$

Φ_i : azimuth of part. in lab.

w_i : weight: p_T ,

opposite sign for/backward
rapidity in case $n = 1$

S. Voloshin and Y. Zhang,

Z. Phys. C 70, 665 (1996)

The event plane resolution depends on:

- the magnitude of the flow of order n
→ **beam energy** (E_{beam}) and **impact parameter** (b) dependant
- the event multiplicity → idem
- the **detector acceptance** and **azimuthal symmetry**